

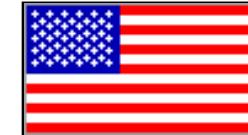
# Elliptic flow, eccentricity and eccentricity fluctuations

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for the  collaboration

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**XXXVI International Symposium on Multiparticle Dynamics**  
**September 02-08, 2006, Paraty, Rio de Janeiro, Brazil**

# PHOBOS collaboration

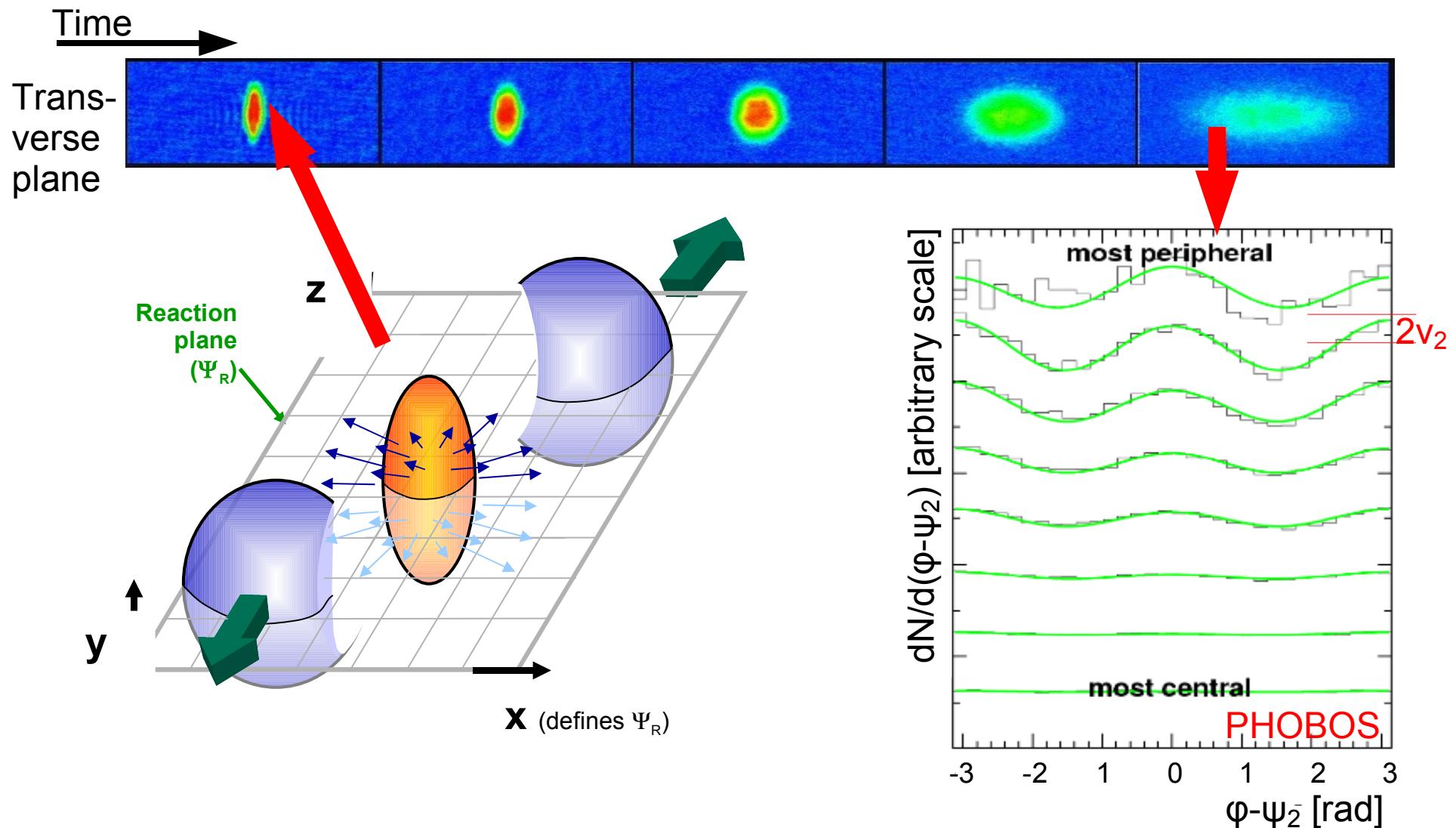


Burak Alver, Birger Back, Mark Baker, Maarten Ballintijn, Donald Barton, Russell Betts, Richard Bindel, Wit Busza (Spokesperson), Zhengwei Chai, Vasundhara Chetluru, Edmundo García, Tomasz Gburek, Kristjan Gulbrandsen, Clive Halliwell, Joshua Hamblen, Ian Harnarine, Conor Henderson, David Hofman, Richard Hollis, Roman Hołyński, Burt Holzman, Aneta Iordanova, Jay Kane, Piotr Kulinich, Chia Ming Kuo, Wei Li, Willis Lin, Constantin Loizides, Steven Manly, Alice Mignerey, Gerrit van Nieuwenhuizen, Rachid Nouicer, Andrzej Olszewski, Robert Pak, Corey Reed, Eric Richardson, Christof Roland, Gunther Roland, Joe Sagerer, Iouri Sedykh, Chadd Smith, Maciej Stankiewicz, Peter Steinberg, George Stephans, Andrei Sukhanov, Artur Szostak, Marguerite Belt Tonjes, Adam Trzupek, Sergei Vaurynovich, Robin Verdier, Gábor Veres, Peter Walters, Edward Wenger, Donald Willhelm, Frank Wolfs, Barbara Wosiek, Krzysztof Woźniak, Shaun Wyngaardt, Bolek Wysłouch

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MASSACHUSETTS INSTITUTE OF TECHNOLOGY  
UNIVERSITY OF ILLINOIS AT CHICAGO  
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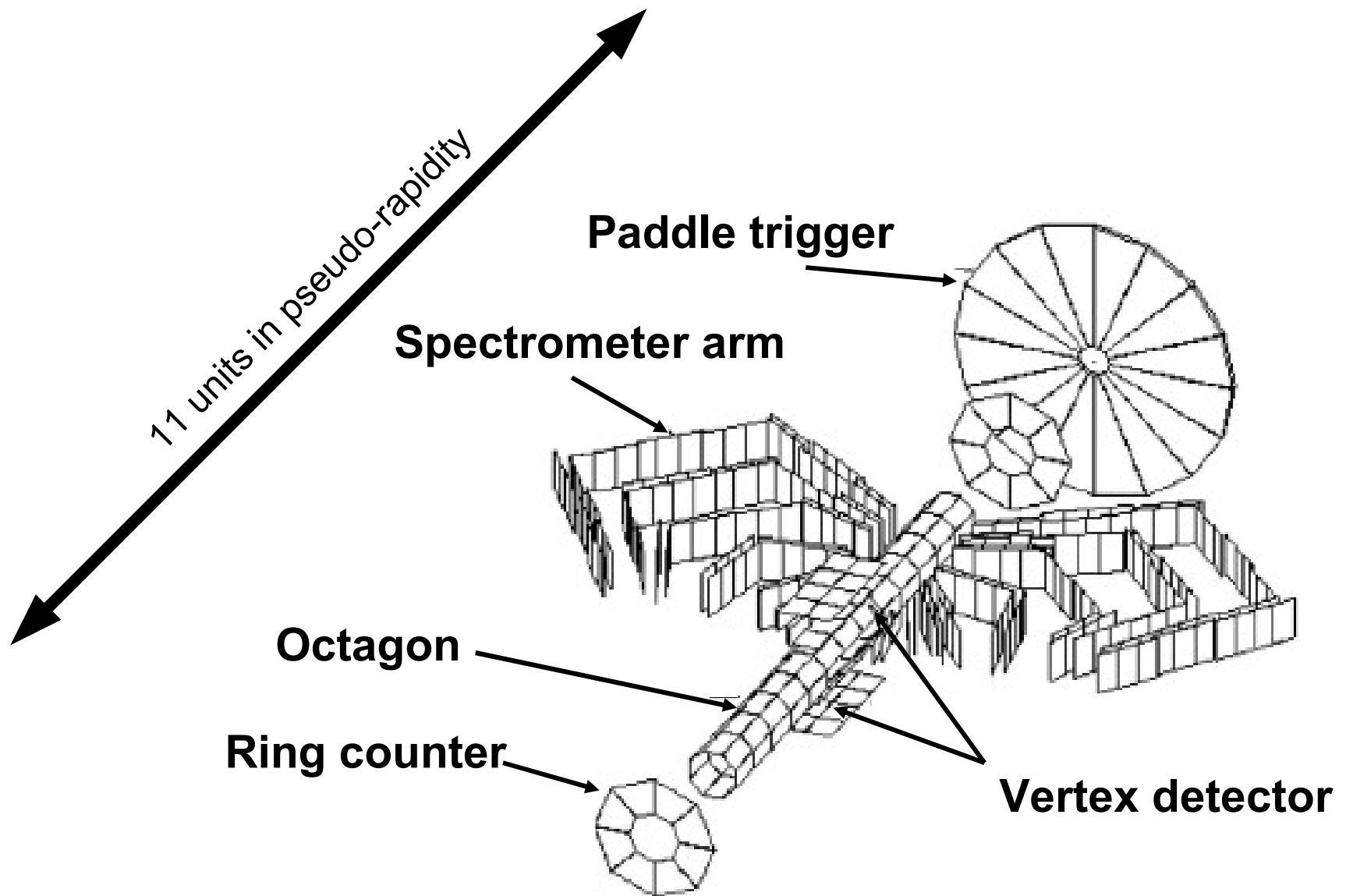
# Elliptic flow



Non-central collision:  
Initial state eccentricity

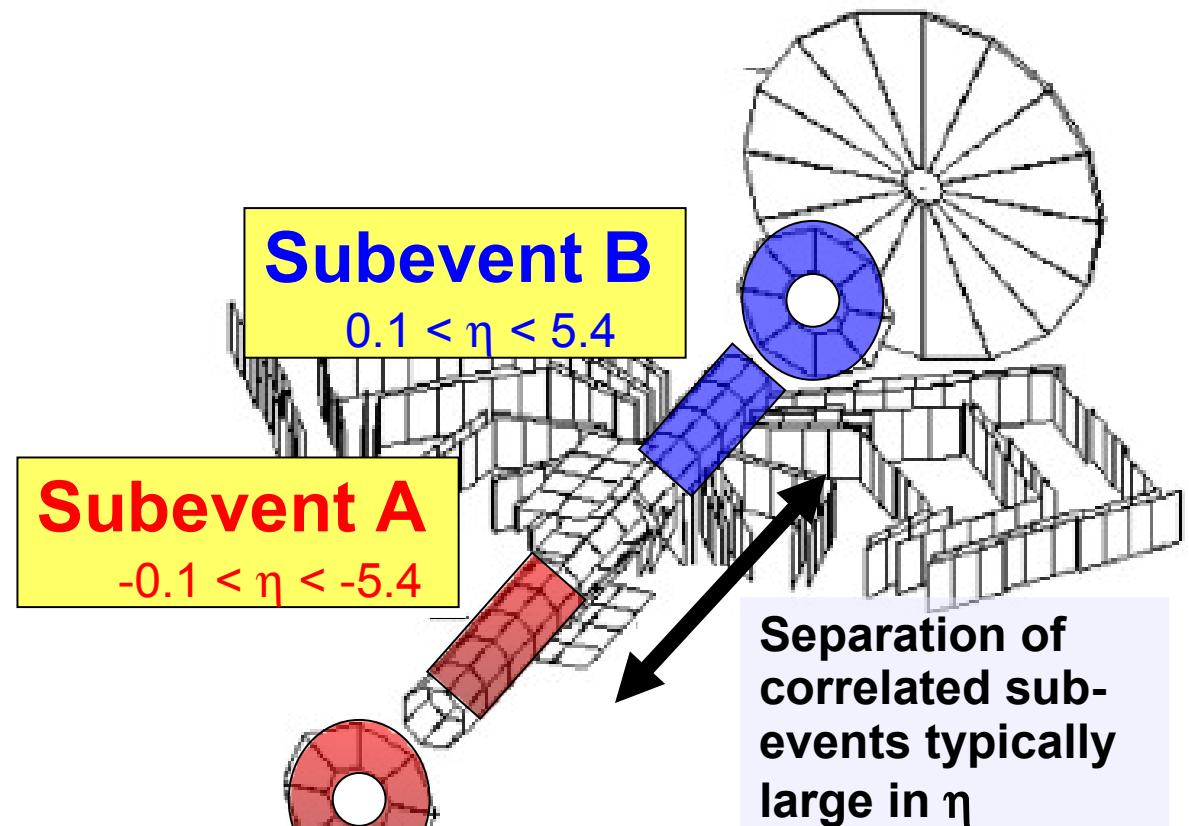
Momentum space anisotropy  
 $v_2 = \langle \cos(2\phi - 2\Psi_R) \rangle$

# Flow measurement in PHOBOS



# Flow measurement in PHOBOS

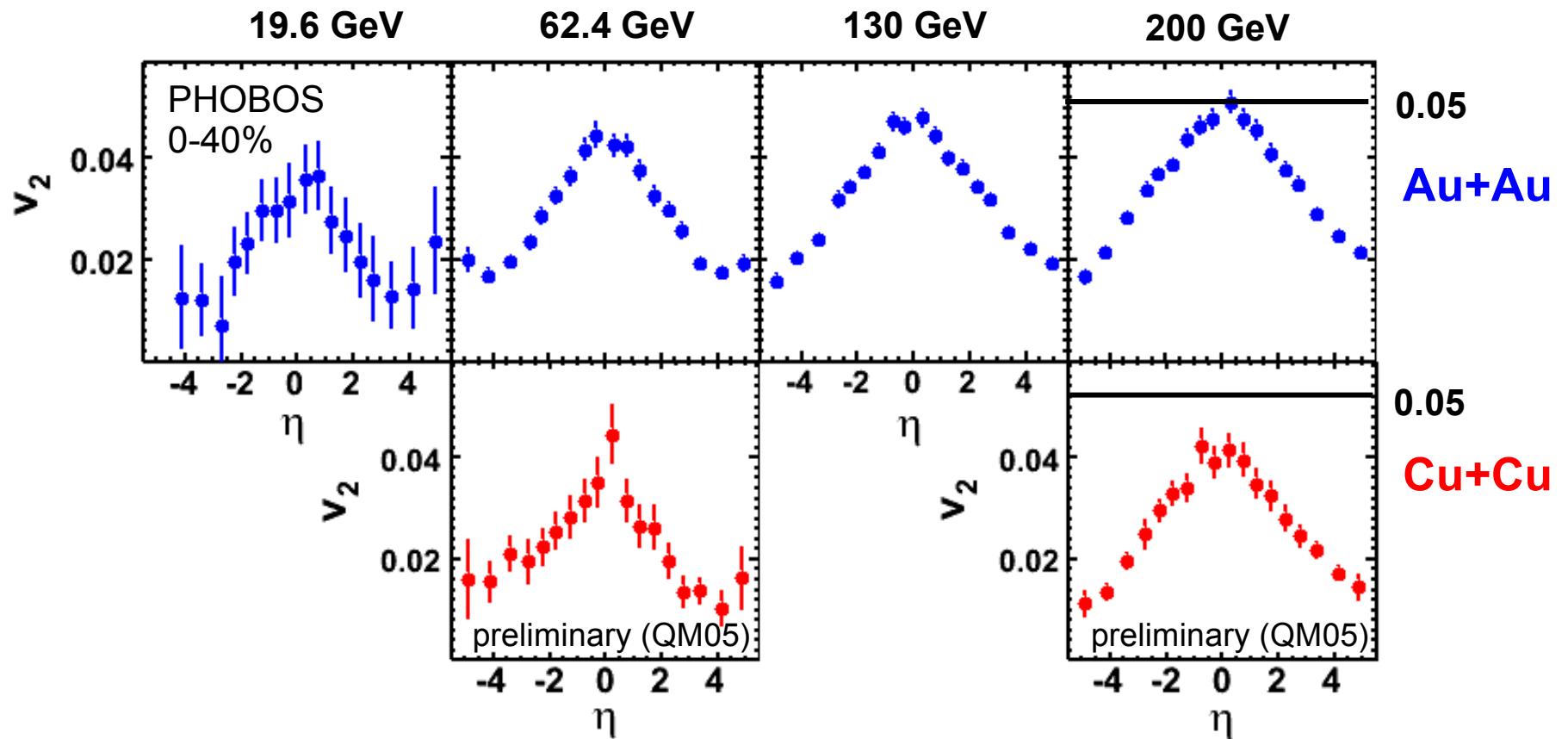
- Reaction-plane / Subevent technique
  - Correlate reaction plane determined from azimuthal pattern of hits in one part of the detector with information from other parts of the detector



$$\tan(2\psi_A) = \frac{\langle \sin(2\phi) \rangle_A}{\langle \cos(2\phi) \rangle_A}$$
$$v_2^{obs} = \langle \cos(2\phi - 2\psi_A) \rangle_B$$
$$v_2 = \frac{\langle v_2^{obs} \rangle_{events}}{\sqrt{\langle \cos(2\psi_A - 2\psi_B) \rangle_{events}}}$$

A.Poskanzer, S.Voloshin,  
nucl-ex/9805001

# Elliptic flow ( $v_2$ )



Cu+Cu about 20% lower than Au+Au

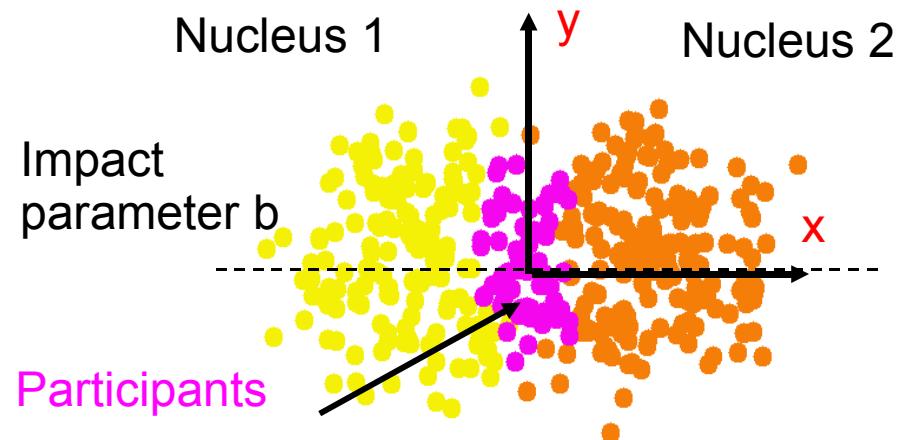
Au+Au: PRL 94 122303 (2005)

Cu+Cu: nucl-ex/0510031 (preliminary)

# Relating system size and eccentricity

- Glauber Monte Carlo
  - Radial distribution of nucleons (in nucleus) drawn from Wood-Saxon distribution
  - Isotropic angular distribution
  - Separate by impact parameter
  - Nucleons travel on straight-line paths and interact inelastically when

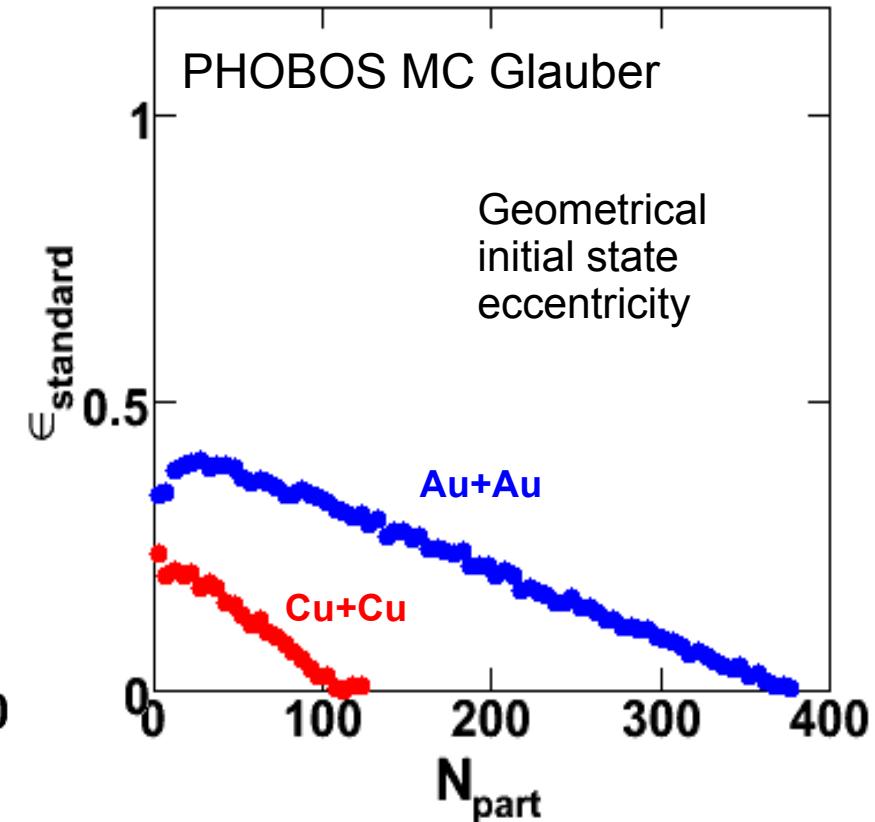
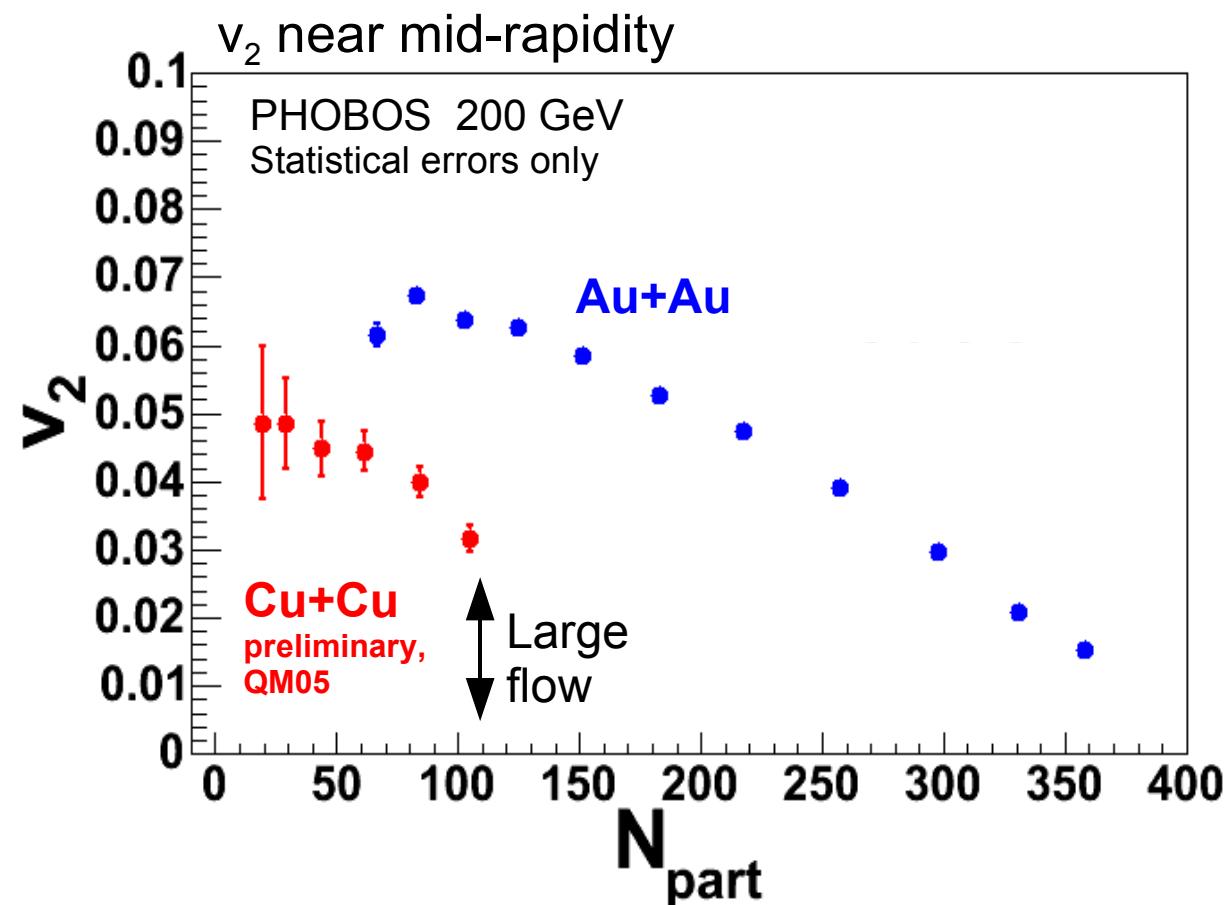
$$d = \sqrt{(x_1 - x_2)^2 + (y_1 - y_2)^2} < \sqrt{\sigma_{NN} / \pi}$$



- Centrality of collision
  - **#Participants**
    - Nucleons that interact at least once
  - Related to cross section and impact parameter range
- **Eccentricity** of collision zone
  - Given by participants position distributions

$$\text{Eccentricity: } \epsilon_{std} = \frac{\sigma_y^2 - \sigma_x^2}{\sigma_y^2 + \sigma_x^2}$$

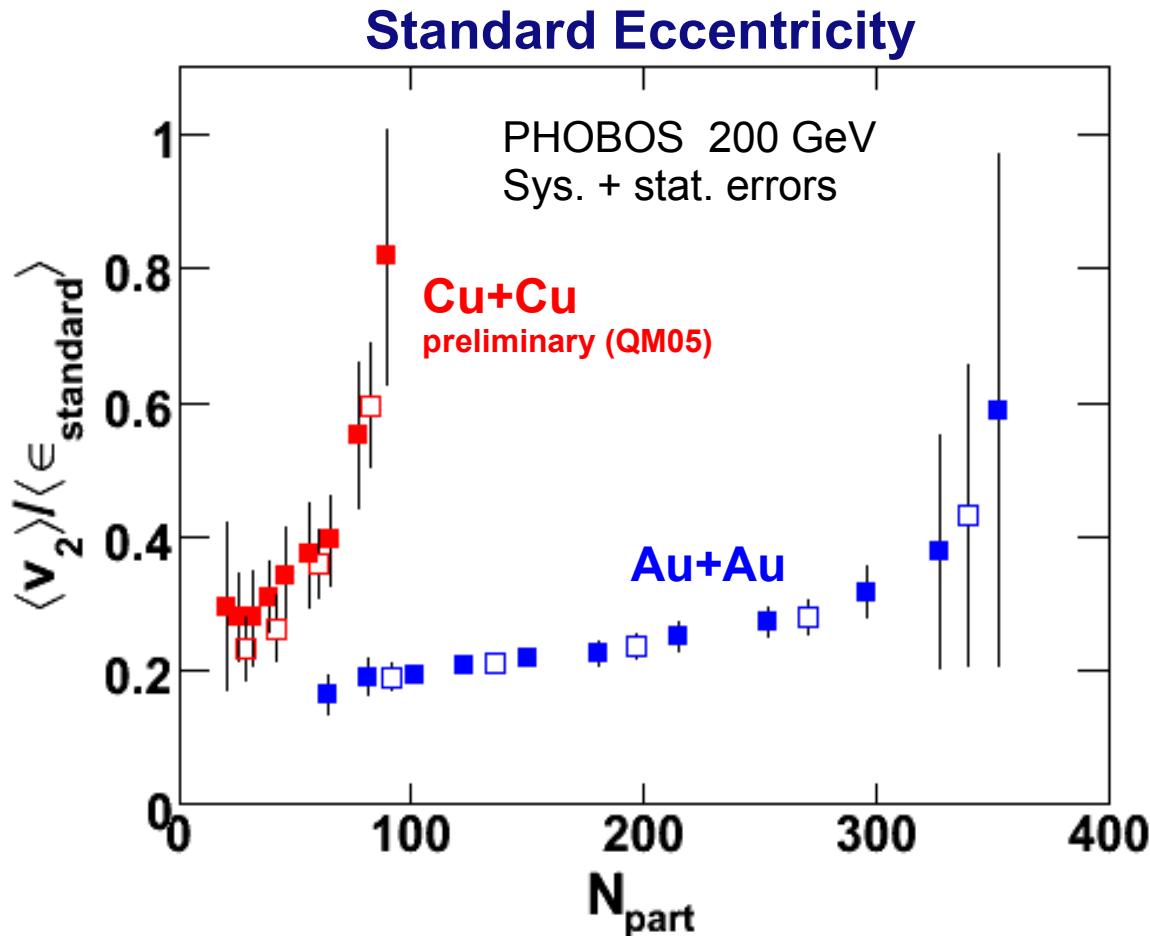
# Elliptic flow in Cu+Cu and Au+Au



Au+Au: PRC 72, 051901 (2005)

Cu+Cu: prel. QM05, nucl-ex/0510042

# Scaled elliptic flow vs $N_{\text{part}}$



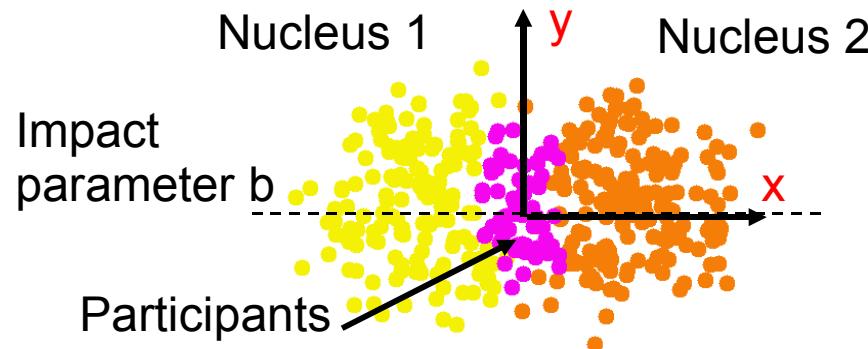
No agreement between Cu+Cu and Au+Au scaled by  $\epsilon_{\text{standard}}$

Au+Au: PRC 72, 051901 (2005)

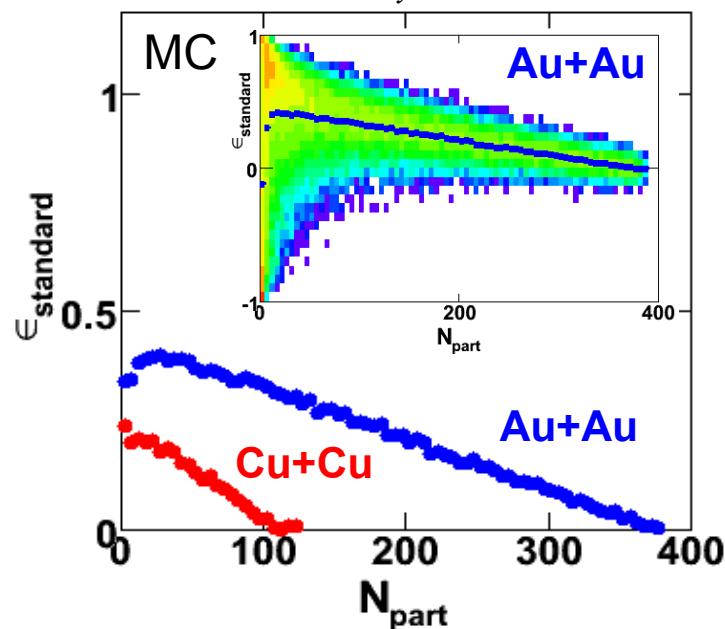
Cu+Cu: prel. QM05, nucl-ex/0510042

# Standard eccentricity calculation

## Standard Eccentricity

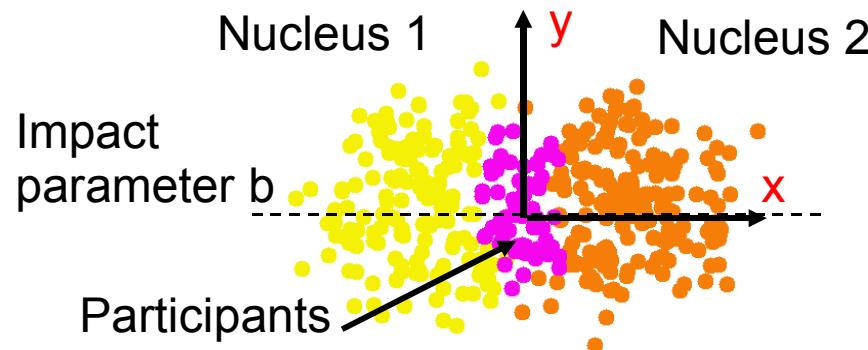


$$\epsilon_{std} = \frac{\sigma_y^2 - \sigma_x^2}{\sigma_y^2 + \sigma_x^2}$$

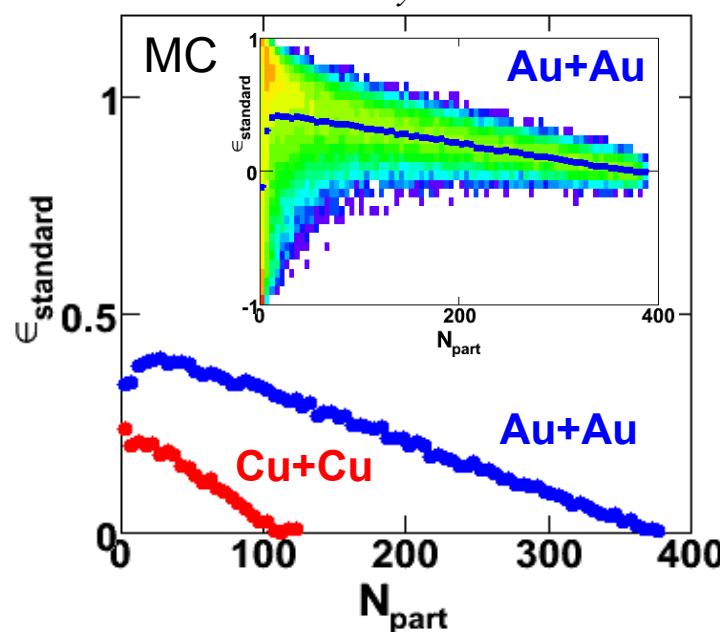


# Participant eccentricity calculation

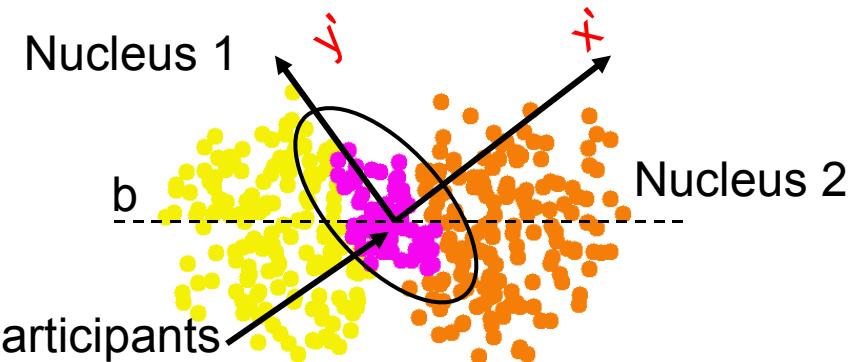
## Standard Eccentricity



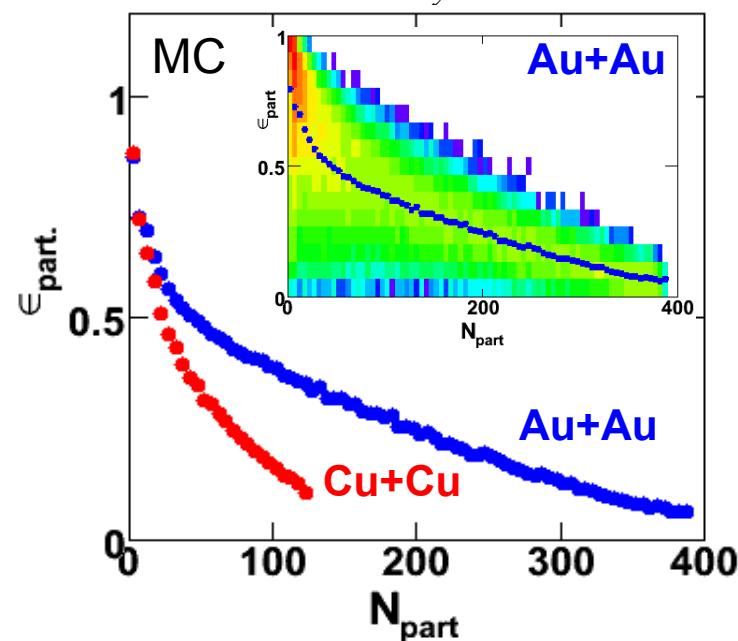
$$\epsilon_{std} = \frac{\sigma_y^2 - \sigma_x^2}{\sigma_y^2 + \sigma_x^2}$$



## Participant Eccentricity



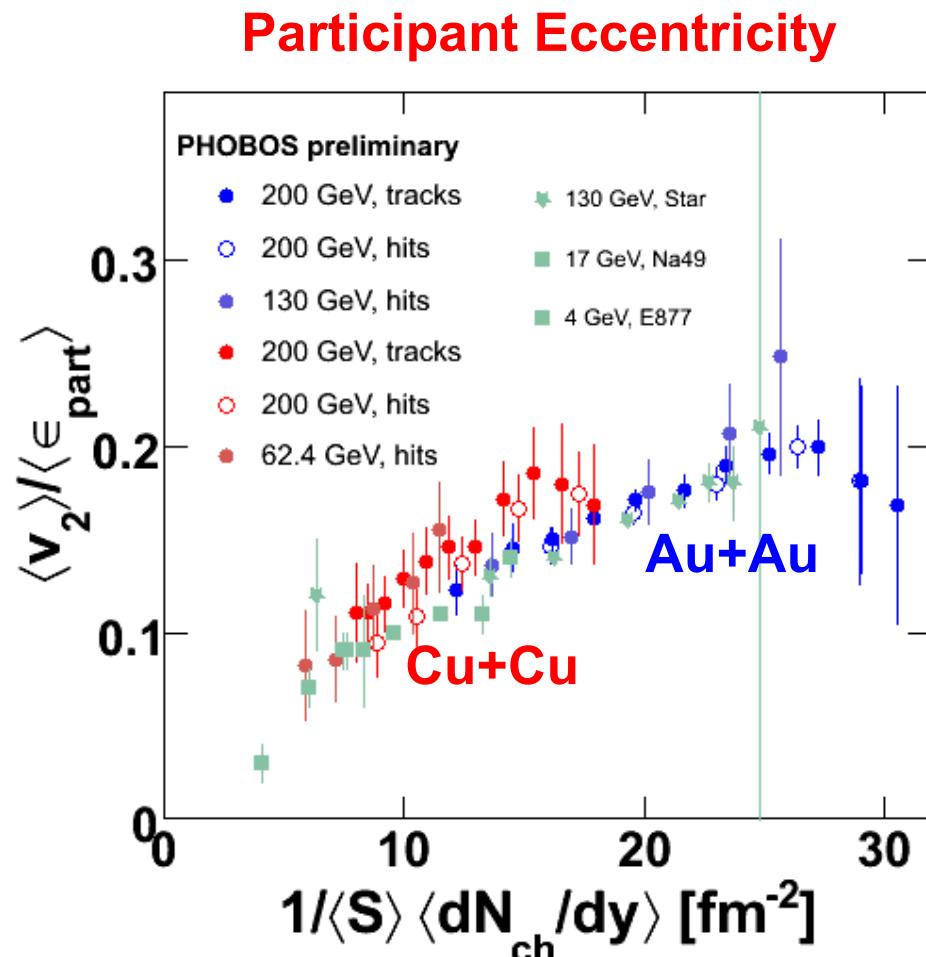
$$\epsilon_{part} = \frac{\sqrt{(\sigma_y^2 - \sigma_x^2)^2 - 4\sigma_{xy}^2}}{\sigma_y^2 + \sigma_x^2}$$



# Participant eccentricity scaling

- Caution: We used  $\epsilon_{\text{participant}}$  for PHOBOS data. Important for Cu-Cu, less critical for Au-Au.
- Scale  $v_2(\eta)$  to  $\sim v_2(y)$  (10% lower)
- Scale  $dN/d\eta$  to be  $\sim dN/dy$  (15% higher)
- STAR and AGS Au+Au and CERN Pb+Pb results are not scaled by  $\epsilon_{\text{participant}}$
- S is overlap area (MC Glauber)

STAR, PRC 66 034904 (2002)  
Voloshin, Poskanzer, PLB 474 27 (2000)  
Heiselberg, Levy, PRC 59 2716, (1999)  
Au+Au: PRC 72, 051901 (2005)  
Cu+Cu: prel. QM05, nucl-ex/0510042

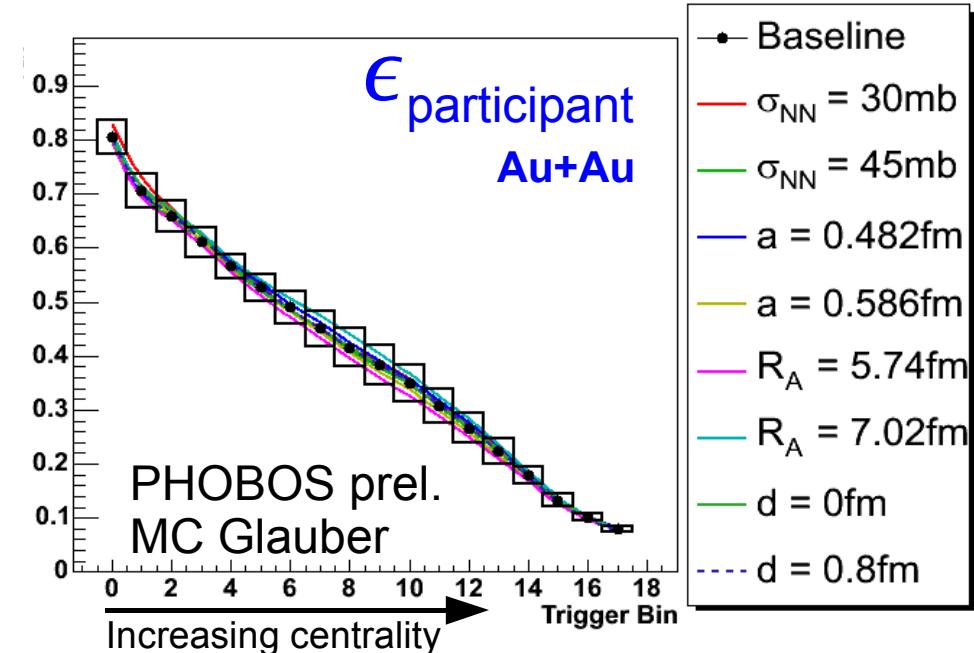
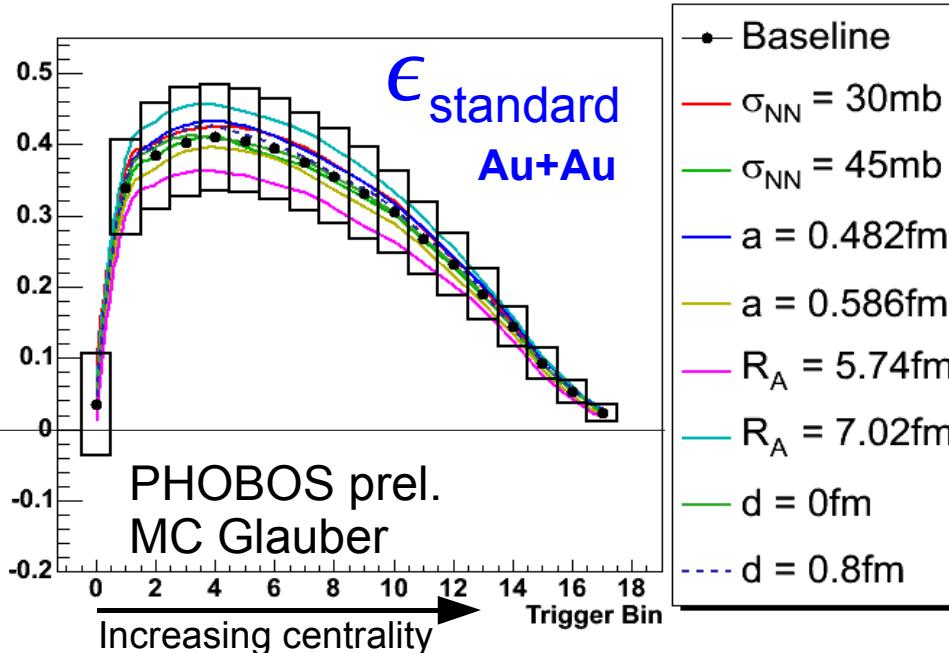


Approximate scaling  
between Cu+Cu and Au+Au

# Since QM 2005

- Examine properties of participant eccentricity
  - Study robustness wrt to geometry parameters
  - Connection to  $v_2$  measurements and hydro
    - Work in collaboration with Uli Heinz
- Measure higher moments in  $v_2$

# Robustness with geometry variables



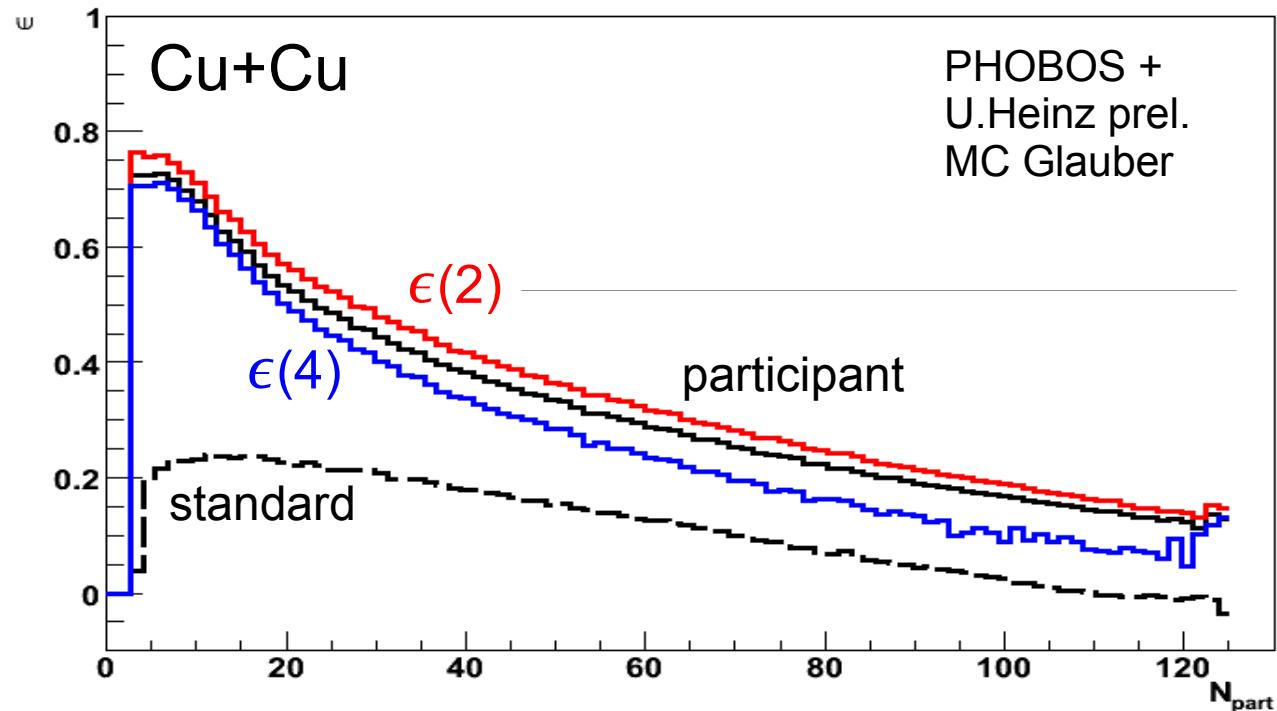
- Variation of
  - Nucleon-nucleon cross section (30-45mb)
  - Nuclear radius ( $\pm 10\%$  from the nominal value)
  - Skin depth (0.482-0.586fm)
  - Minimum separation distance between nucleons ( $d=0-0.8\text{fm}$ )

$$\rho(r) = \frac{\rho_0}{1 + \exp((r-R)/a)}$$

$\epsilon_{\text{participant}}$  even slightly more robust than  $\epsilon_{\text{standard}}$

# Moments of eccentricity

- If one measures  $v_2$  with two-particle correlations  
 $v_2(2) = \sqrt{\langle \cos 2(\phi_1 - \phi_2) \rangle} = \sqrt{\langle v_2^2 \rangle}$
- If  $v_2$  fluctuates prop. to  $\epsilon_{\text{part}}$   
 $\epsilon_{\text{part}} = \epsilon(2) = \sqrt{\langle \epsilon_{\text{part}}^2 \rangle}$   
Scaling with
- For  $\epsilon(4)$  cumulant method  
 $\epsilon(4) = [2\langle \epsilon^2 \rangle^2 - \langle \epsilon^4 \rangle]^{1/4}$   
Scaling with



$\epsilon\{2\}$  rather similar to  $\epsilon_{\text{part}}$ , unlike  $\epsilon\{4\}$   
 that is between  $\epsilon_{\text{std}}$  and  $\epsilon_{\text{part}}$

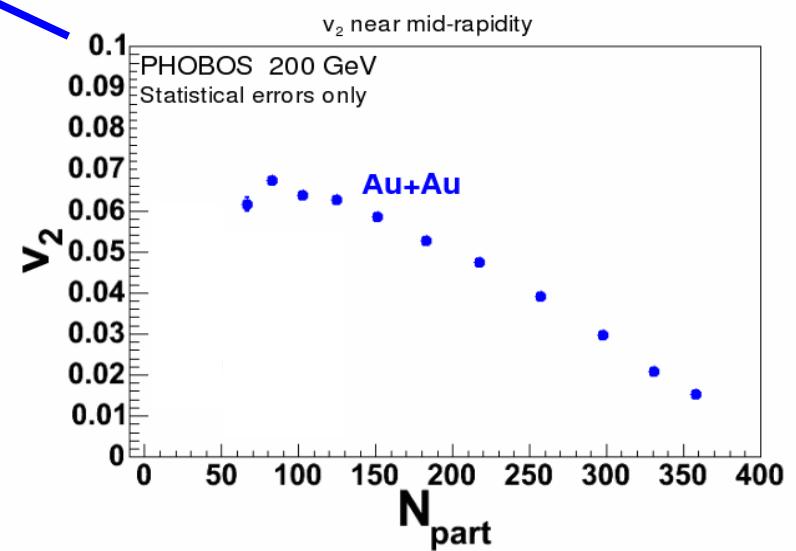
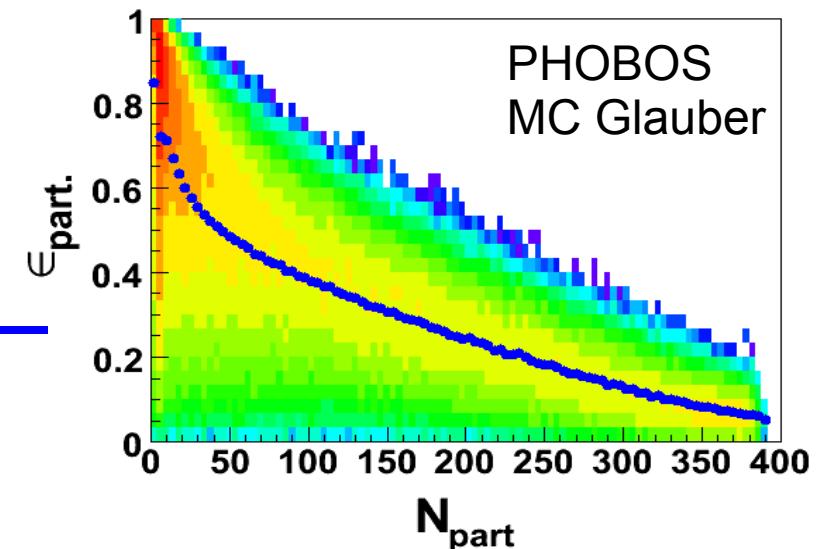
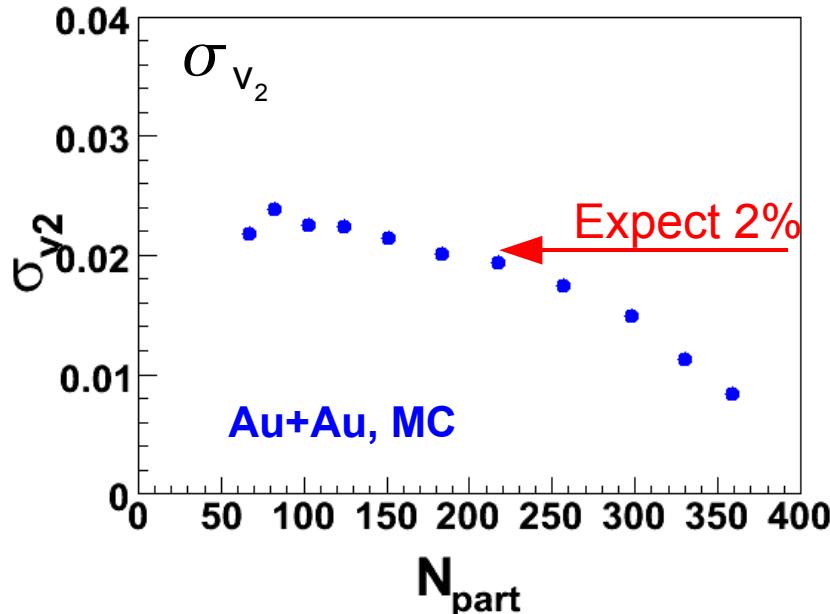
M.Miller, R.Snellings, nucl-ex/0312008  
 R.Bhalerao, J.Ollitrault, nucl-th/0607009

# Expected elliptic flow fluctuations

Assuming  $v_2 \propto \epsilon_{\text{part}}$ , the participant eccentricity model predicts

$$\sigma_{v_2} = \frac{\sigma_{\epsilon_{\text{part}}}}{\epsilon_{\text{part}}} v_2$$

Expected  $\sigma_{v_2}$  from fluctuations in  $\epsilon_{\text{part}}$



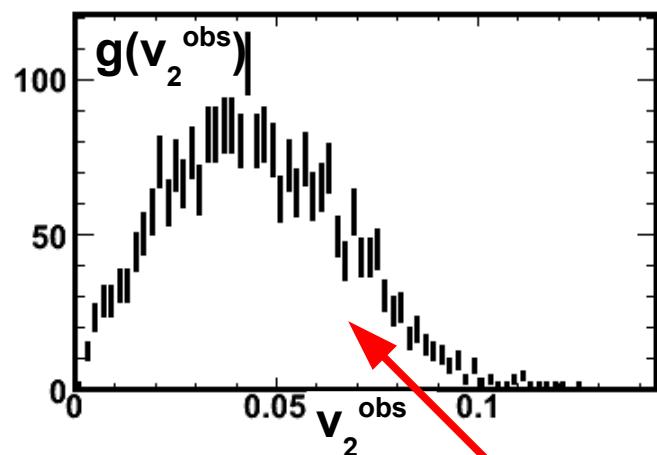
# Measuring elliptic flow fluctuations

- Ongoing analysis on 200GeV Au+Au
  - NO DATA will be shown
  - Overview of the measurement
    - Studies on fully simulated MC events
  - Details are in nucl-ex/0608025

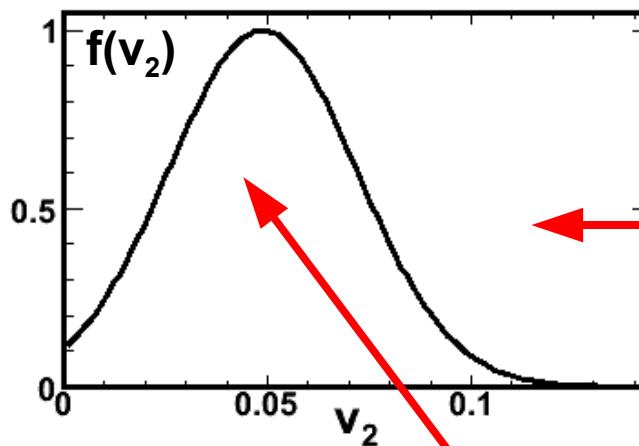
B.Alver et.al. (PHOBOS), nucl-ex/0608025

# Measuring elliptic flow fluctuations

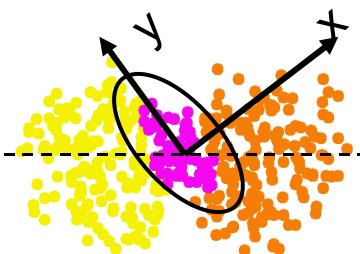
Observed  $v_2$  distribution



True  $v_2$  distribution

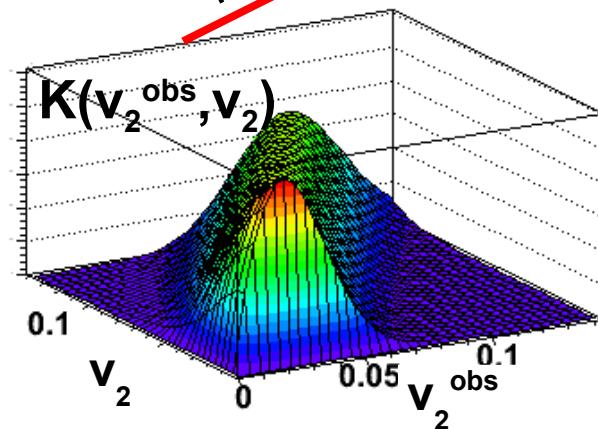


Participant eccentricity

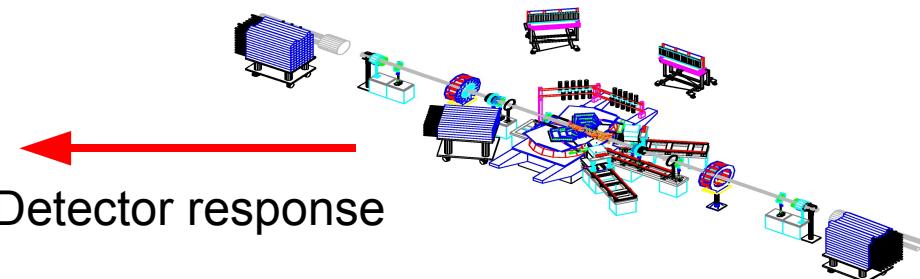


$$g(v_2^{\text{obs}}) = \int_0^\infty K(v_2^{\text{obs}}, v_2) f(v_2) dv_2$$

Kernel

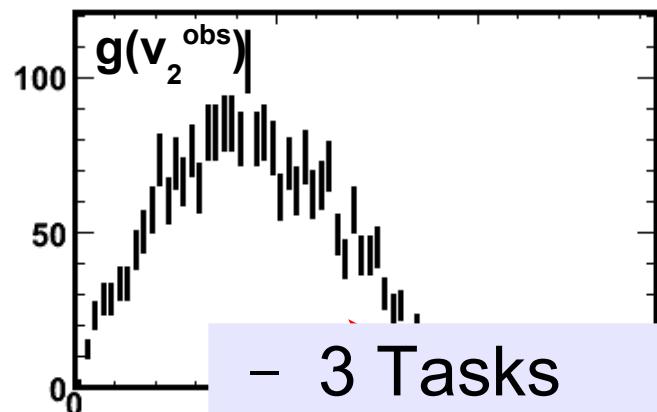


Detector response

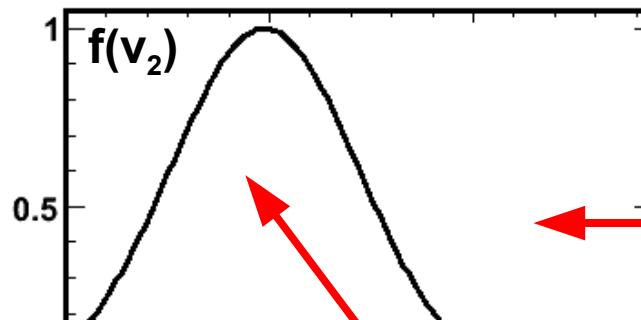


# Measuring elliptic flow fluctuations

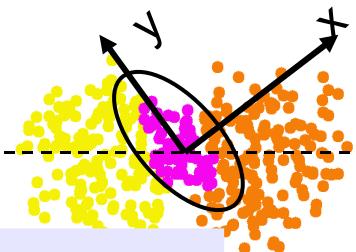
Observed  $v_2$  distribution



True  $v_2$  distribution

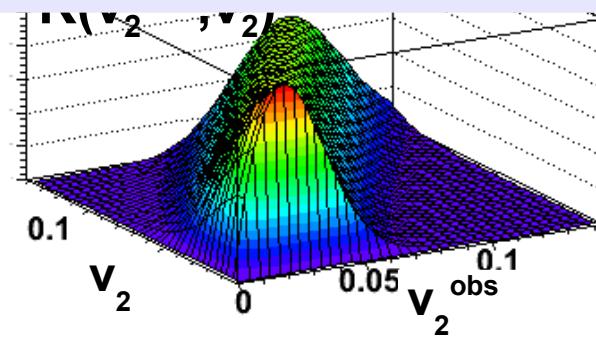


Participant eccentricity

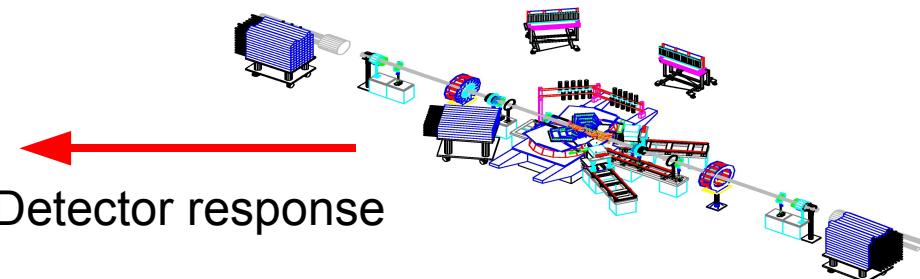


## – 3 Tasks

- Measurement of  $v_2^{obs}$  event-by-event:  
 $g(v_2^{obs})$
- Construction of the kernel:  $K(v_2^{obs}, v_2)$
- Extraction of dynamical fluctuations:  $f(v_2)$



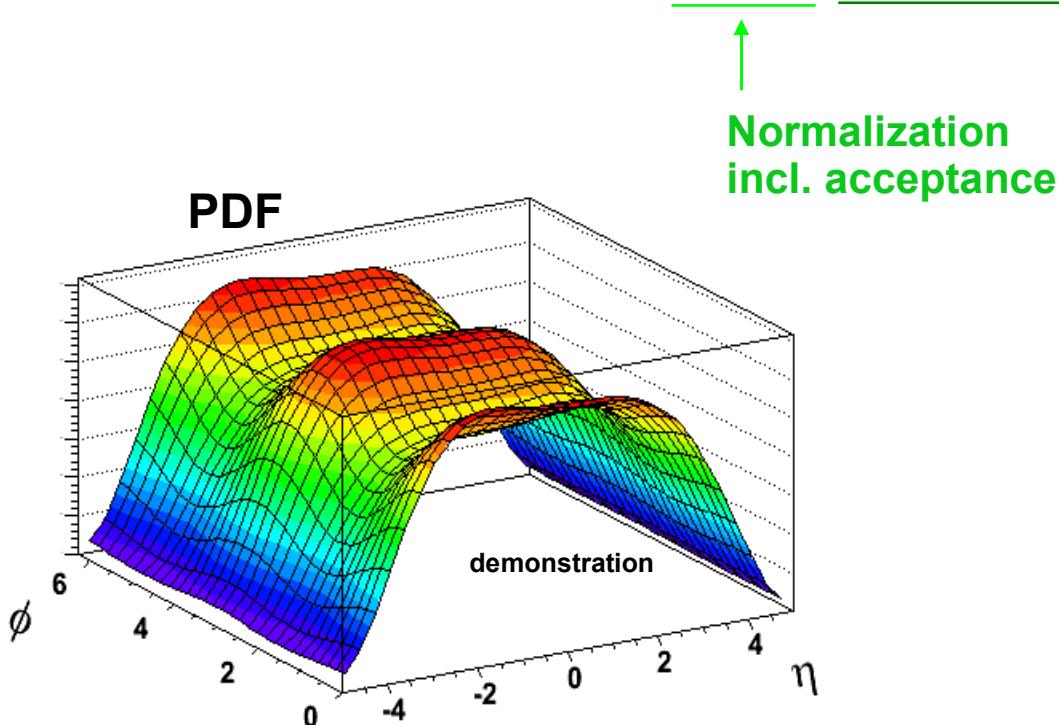
Detector response



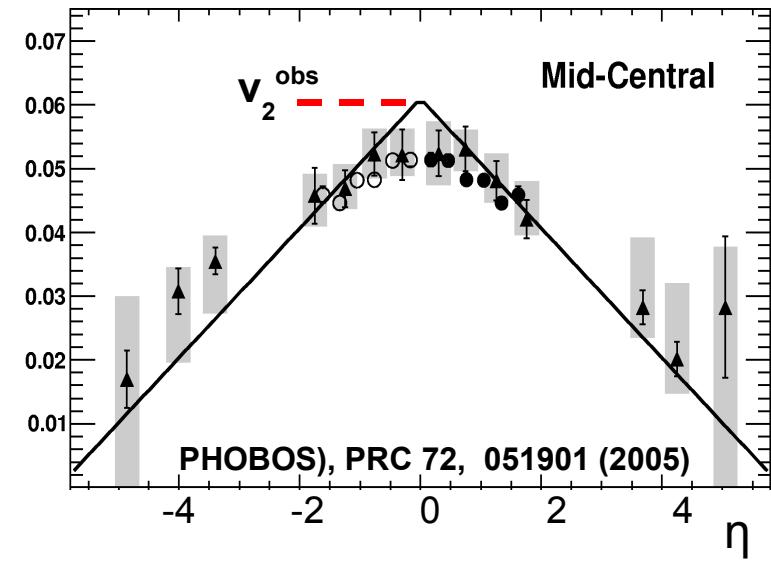
# Event-by-event measurement of $v_2^{\text{obs}}$

- Probability Distribution Function (PDF) for hit positions:

$$P(\eta, \phi; v_2^{\text{obs}}, \phi_0) = p(\eta) [1 + 2v_2(\eta) \cos(2\phi - 2\phi_0)]$$



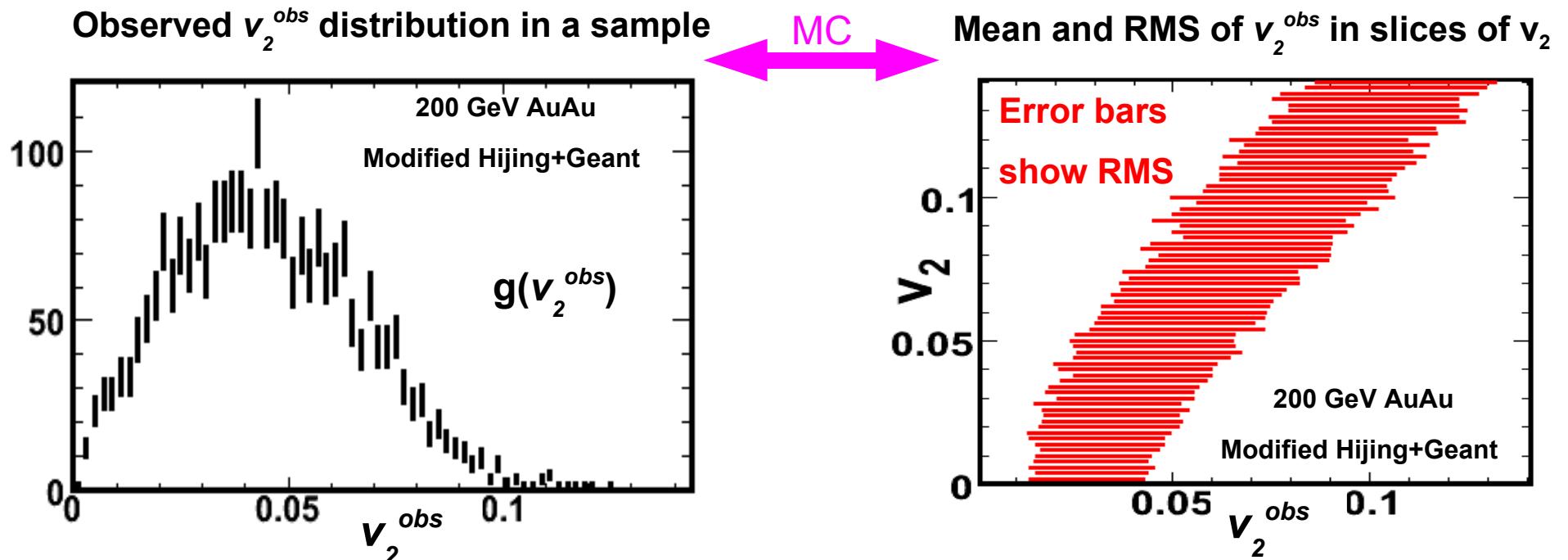
$$v_2(\eta) = v_2^{\text{obs}} \left(1 - \frac{|\eta|}{6}\right)$$



- Maximize the likelihood function

$$L(v_2^{\text{obs}}, \phi_0) = \prod_{i=1}^n P(\eta_i, \phi_i; v_2^{\text{obs}}, \phi_0)$$

# Determining the kernel



- Construction of kernel: “Measure”  $v_2^{obs}$  distribution in bins of  $v_2$  in MC
- 2 small complications
  - Kernel really depends on multiplicity (centrality):  $K(v_2^{obs}, v_2, n)$ 
    - $n$  = number of hits on the detector
    - Determine  $v_2^{obs}$  distribution in bins of  $v_2$  and  $n$ .
  - Deal with low statistics in bins by determining smooth functions

# Extracting dynamical fluctuations

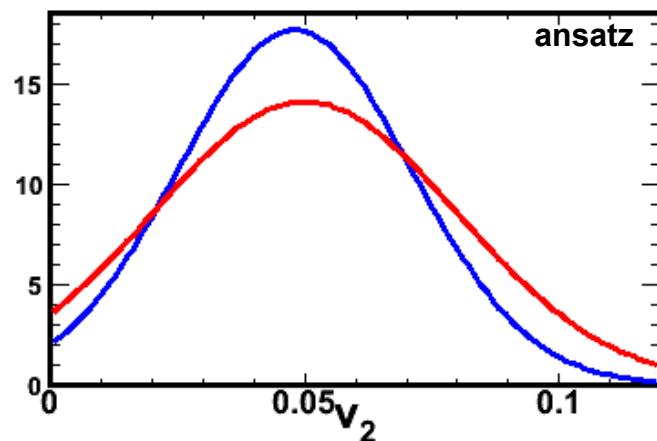
$$g(v_2^{\text{obs}}) = \int_0^\infty K(v_2^{\text{obs}}, v_2) f(v_2) dv_2$$

known

Gaussian Ansatz:

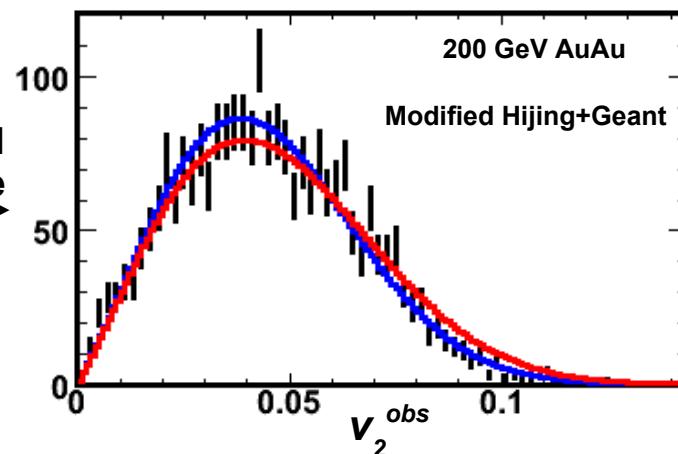
$$f(v_2) = \exp\left[\frac{-(v_2 - \langle v_2 \rangle)^2}{2 \sigma_{v_2}^2}\right]$$

Trials for  $f(v_2)$



Use kernel  
+ integrate

Comparison with sample



Compare expected  $g(v_2^{\text{obs}})$  for Ansatz with measurement

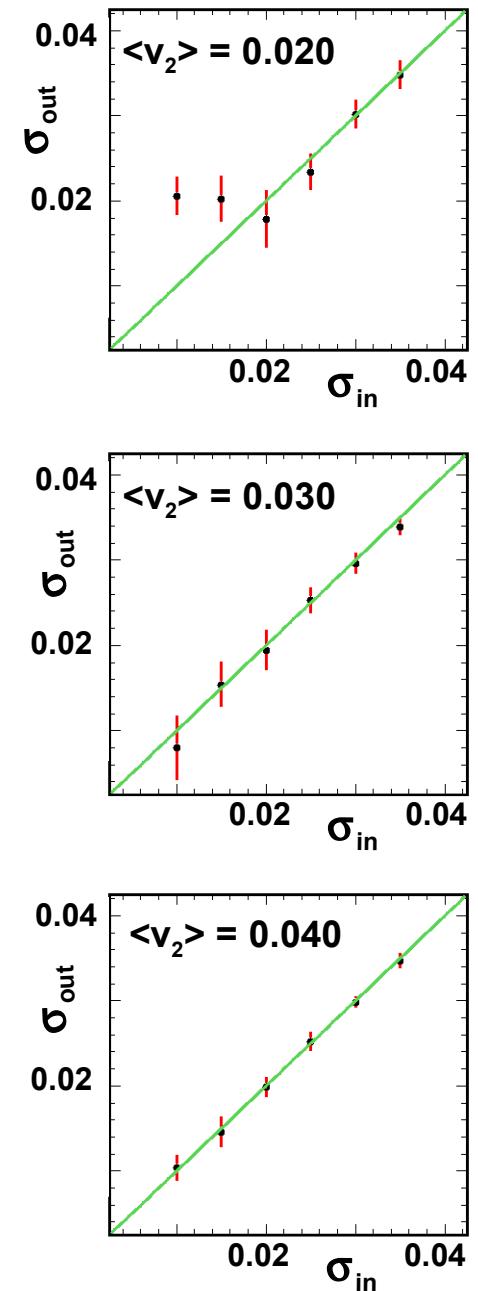
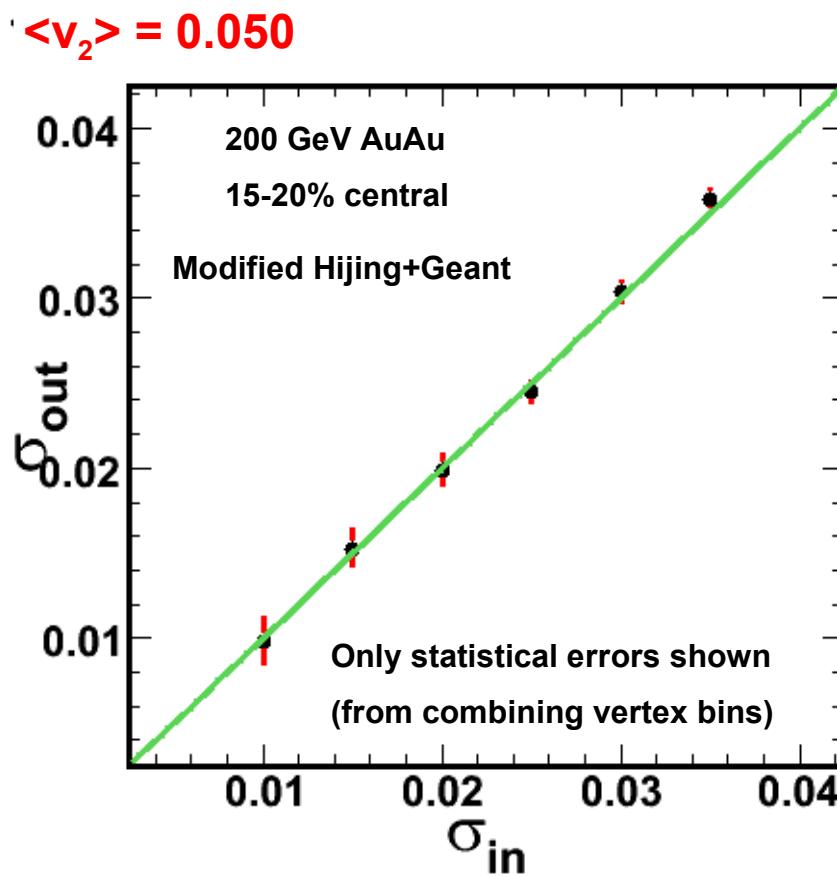
Minimum  $\chi^2 \rightarrow \langle v_2 \rangle$  and  $\sigma_{v_2}$

# Verification of analysis method

- Run analysis on Modified Hijing
  - Shape  $v_2(\eta) = v_2(0) \cdot (1-|\eta|/6)$ 
    - Same as the assumption in our pdf
  - Analysis done in 10 collision vertex bins
    - Final results are averaged
- Kernel
  - 0-40% central events used to construct the kernel
- Input sample
  - 15-20% central events used as sample
  - $v_2(0)$  given by a Gaussian distribution
    - Same as our Ansatz

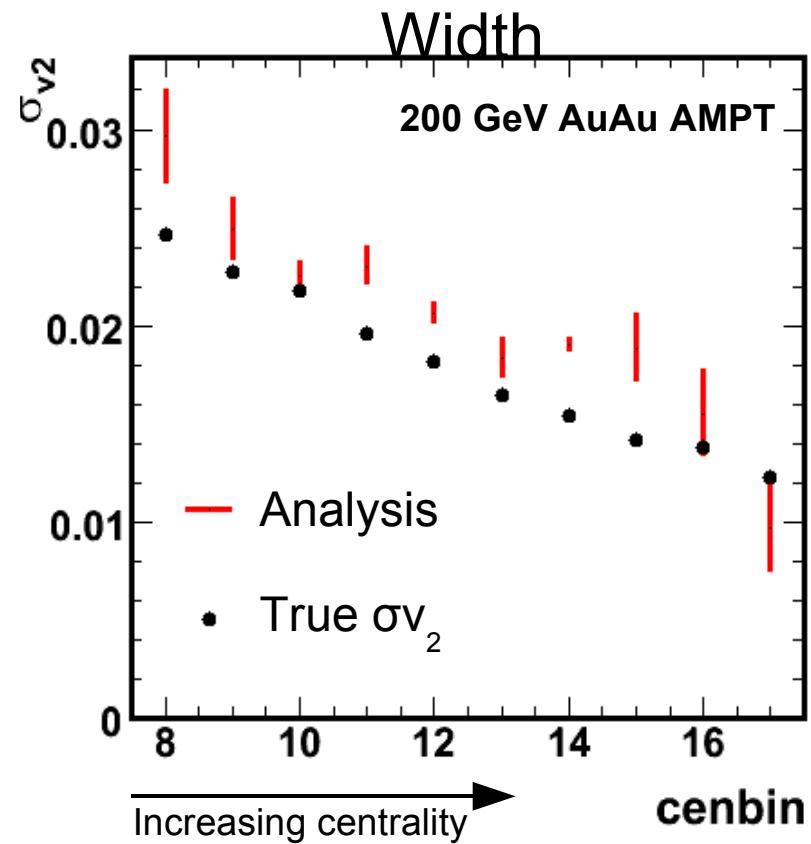
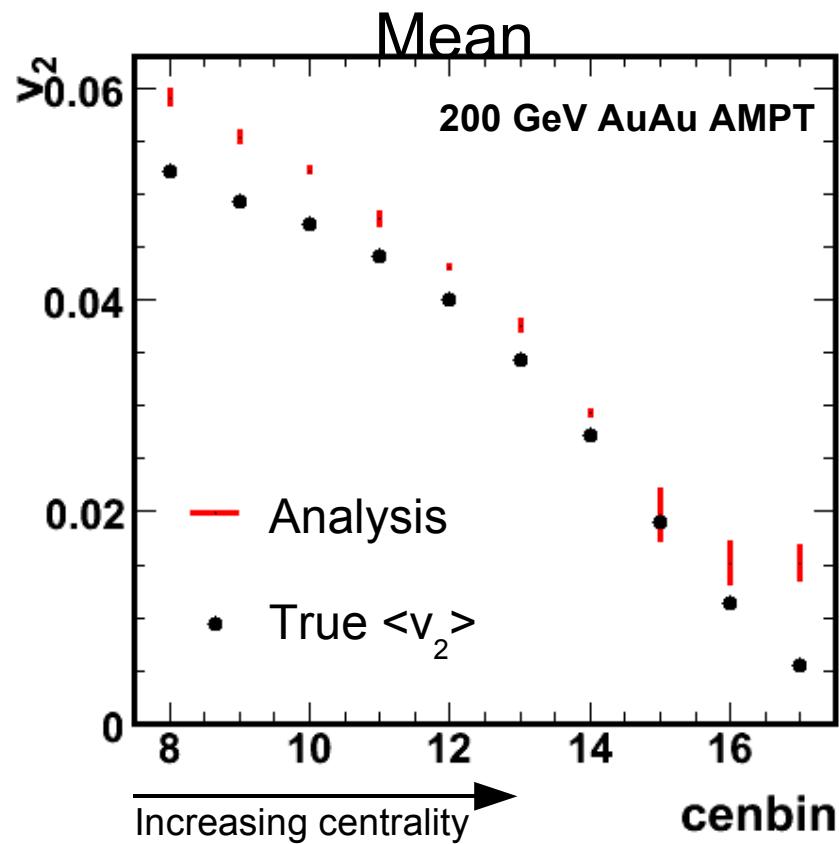
# Verification of analysis method

- Run analysis on Modified Hijing
  - Input fluctuations are successfully reconstructed



# Systematic verification with AMPT

- Run analysis on AMPT
  - Kernel is constructed from Hijing as before
  - Input fluctuations are successfully reconstructed

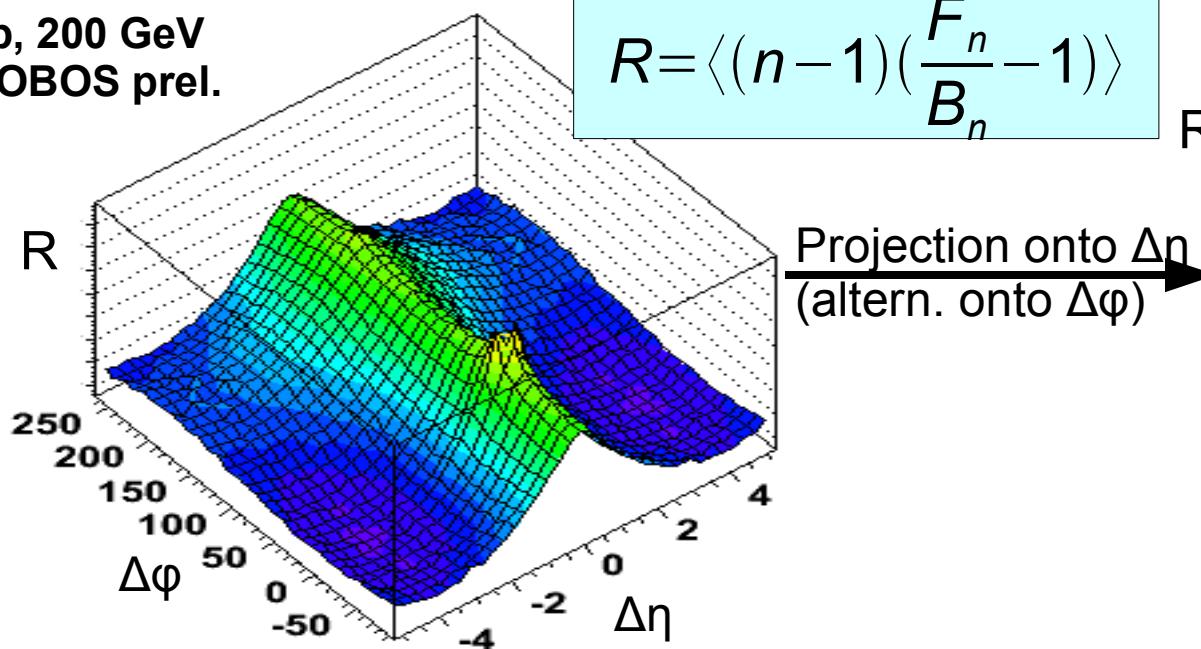


# Non-flow contributions to flow fluctuations

- Non-flow correlations mimic dynamical fluctuations and will contribute to the width of the  $v_2$  distribution
  - The resolution of our method depends on the kernel
    - Modified Hijing: particle multiplicity defines the resolution
    - Data (AMPT): clusters flow and therefore the cluster multiplicity determines the resolution
  - The fluctuations we measure are real (present at particle level) but might not be the ones we are after
- Kernel could compensate for non-flow effects if they are correctly described by the MC used to construct it
  - Construct and tune MC on data
  - Two-particle correlation measurements can be used as input to disentangle the different contributions

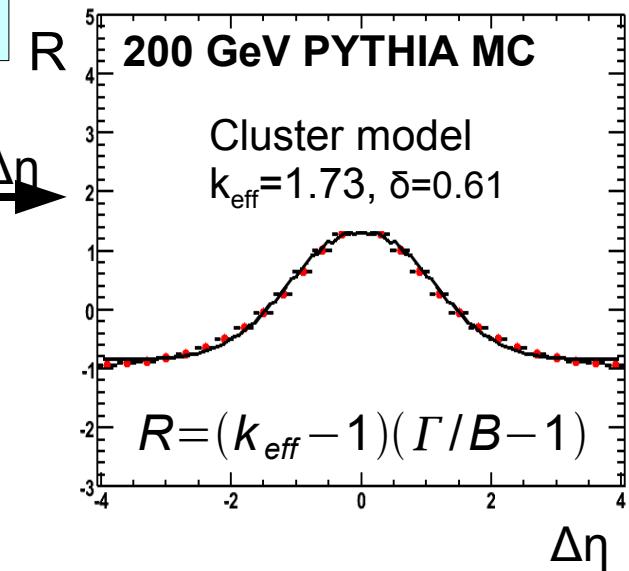
# Two-particle angular correlations in p+p

p+p, 200 GeV  
PHOBOS prel.



$$R = \langle (n-1) \left( \frac{F_n}{B_n} - 1 \right) \rangle$$

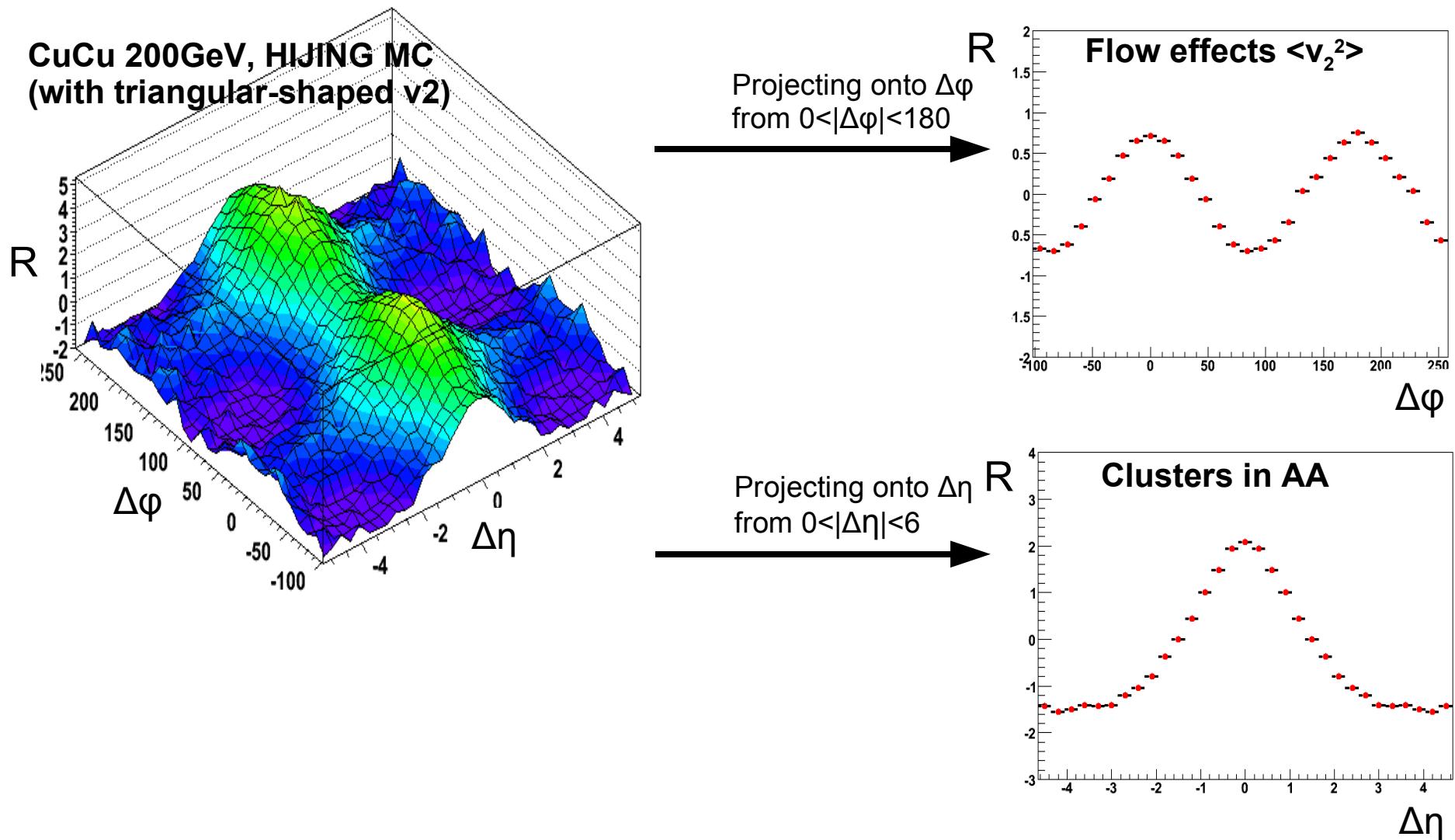
Projection onto  $\Delta\eta$   
(altern. onto  $\Delta\phi$ )



- Construction of  $R$ , event-by-event, weighted by event multiplicity
  - Full  $\phi$  and large  $|\eta| \leq 3$  coverage ( $|\eta| \leq 5$  for future studies)
  - Single hit in silicon layer instead of particle information
    - Need special care for secondary contamination
  - Study soft physics (No trigger particle)
    - Clusters

Will be published soon

# Two-particle angular correlations in A+A



Comprehensive study of two-particle correlations in p+p, d+A and A+A will help distangle different effects in HI systems

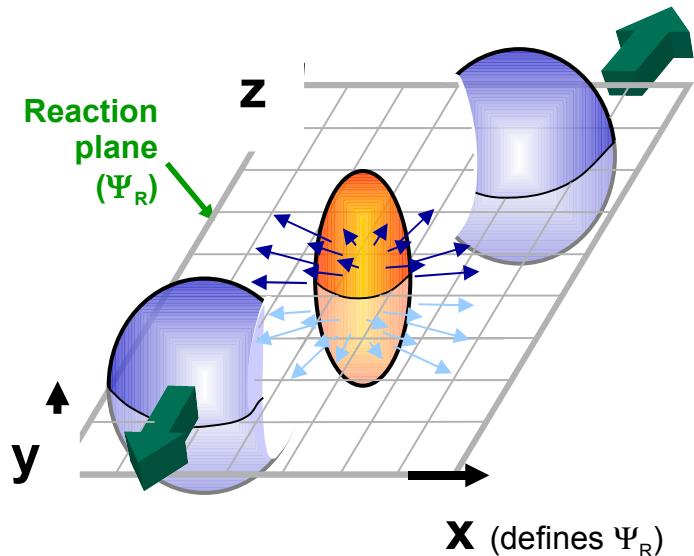
# Summary and Perspectives

- Mean  $v_2$  measurements
  - Large elliptic flow in Cu+Cu compared to Au+Au
- Eccentricity fluctuations
  - Participant eccentricity connects elliptic flow in small system to initial geometry fluctuations
  - Robustness studies wrt to changes in Glauber geometry
  - Work in progress with U.Heinz to investigate/check further properties
- Elliptic flow fluctuations
  - A method to measure elliptic flow fluctuations has been developed
  - Fluctuations in MC simulations are successfully reconstructed
  - Study systematic uncertainties due to MC/DATA differences
- Two-particle angular correlations
  - Properties of clusters in p+p, d+A, and A+A
  - Estimate non-flow contribution to elliptic flow fluctuations
    - Tune MC

# Backup

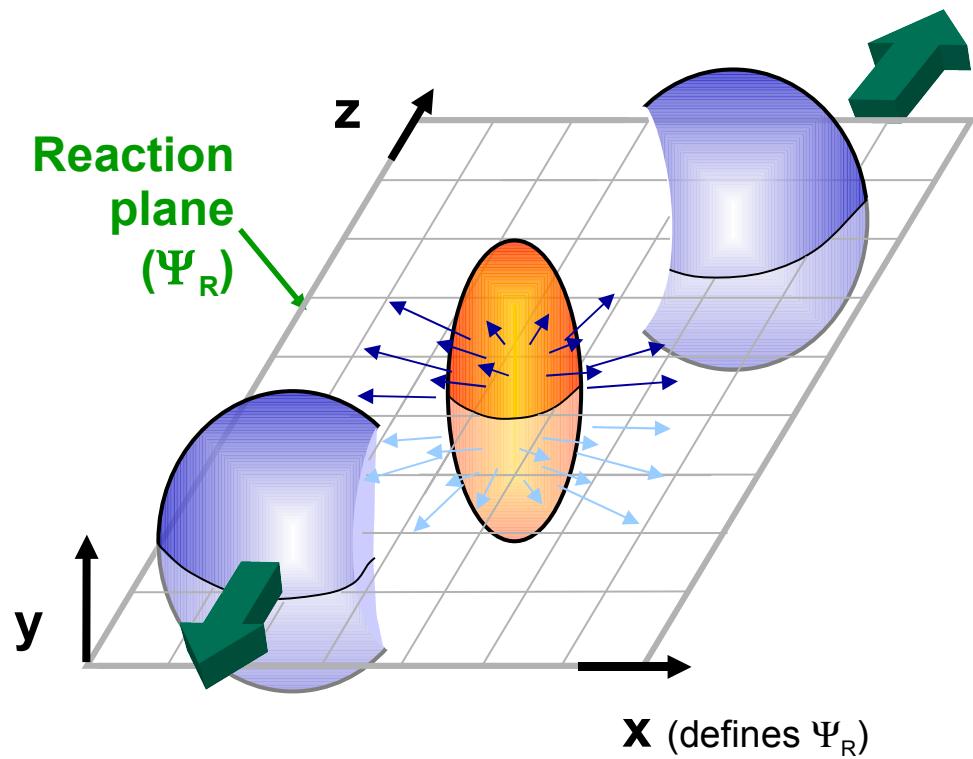
# Anisotropic flow

$$v_n = \langle \cos(n\phi - n\Psi_R) \rangle$$
$$v_n = \langle \cos(n\phi - n\Psi_R) \rangle$$



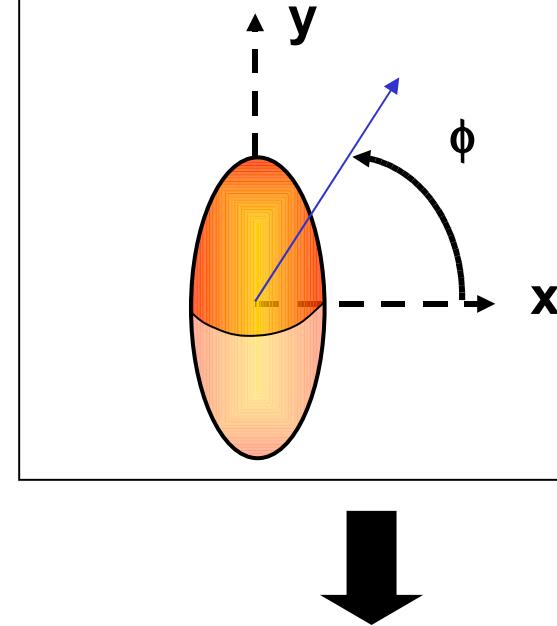
- Anisotropic flow  $\equiv$  azimuthal correlations with reaction plane
  - Consequence of thermalization
  - Final-state reinteractions
  - Typically described by hydrodynamics (but flow does not necessarily imply it)
- Non-flow  $\equiv$  contribution to  $v_n$  from azimuthal correlations between particles (HBT, resonances, jets)

# Direct ( $v_1$ ) and elliptic ( $v_2$ ) flow

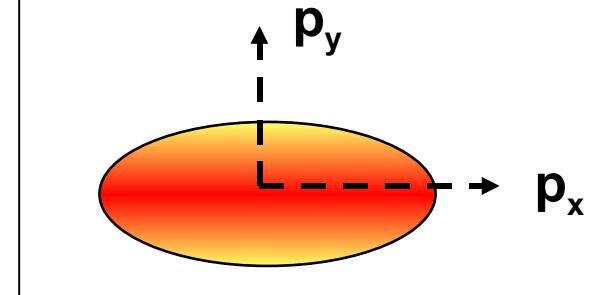


$$\frac{dN}{d(\phi - \Psi_R)} = N_0 (1 + 2v_1 \cos(\phi - \Psi_R) + 2v_2 \cos(2(\phi - \Psi_R)) + \dots)$$

Initial spatial anisotropy



Final momentum anisotropy



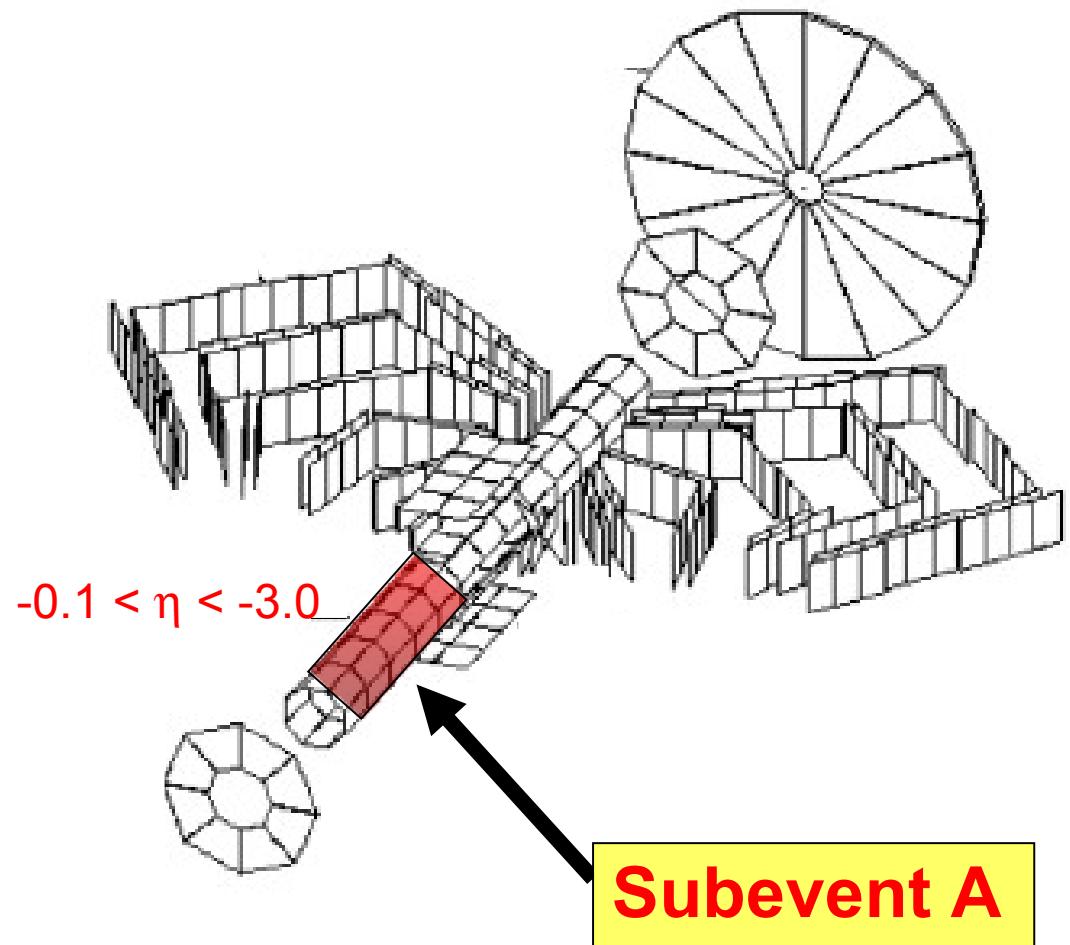
# Flow measurement in PHOBOS

- Reaction-Plane / Subevent technique
  - Correlate reaction plane determined from azimuthal pattern of hits in one part of the detector with information from other parts of the detector
    - Hits
    - Tracks

$$\tan(2\psi_A) = \frac{\langle \sin(2\phi) \rangle_A}{\langle \cos(2\phi) \rangle_A}$$

$$v_2^{obs} = \langle \cos(2\phi - 2\psi_A) \rangle_B$$

$$v_2 = \frac{\langle v_2^{obs} \rangle_{events}}{\sqrt{\langle \cos(\psi_A - \psi_B) \rangle_{events}}}$$



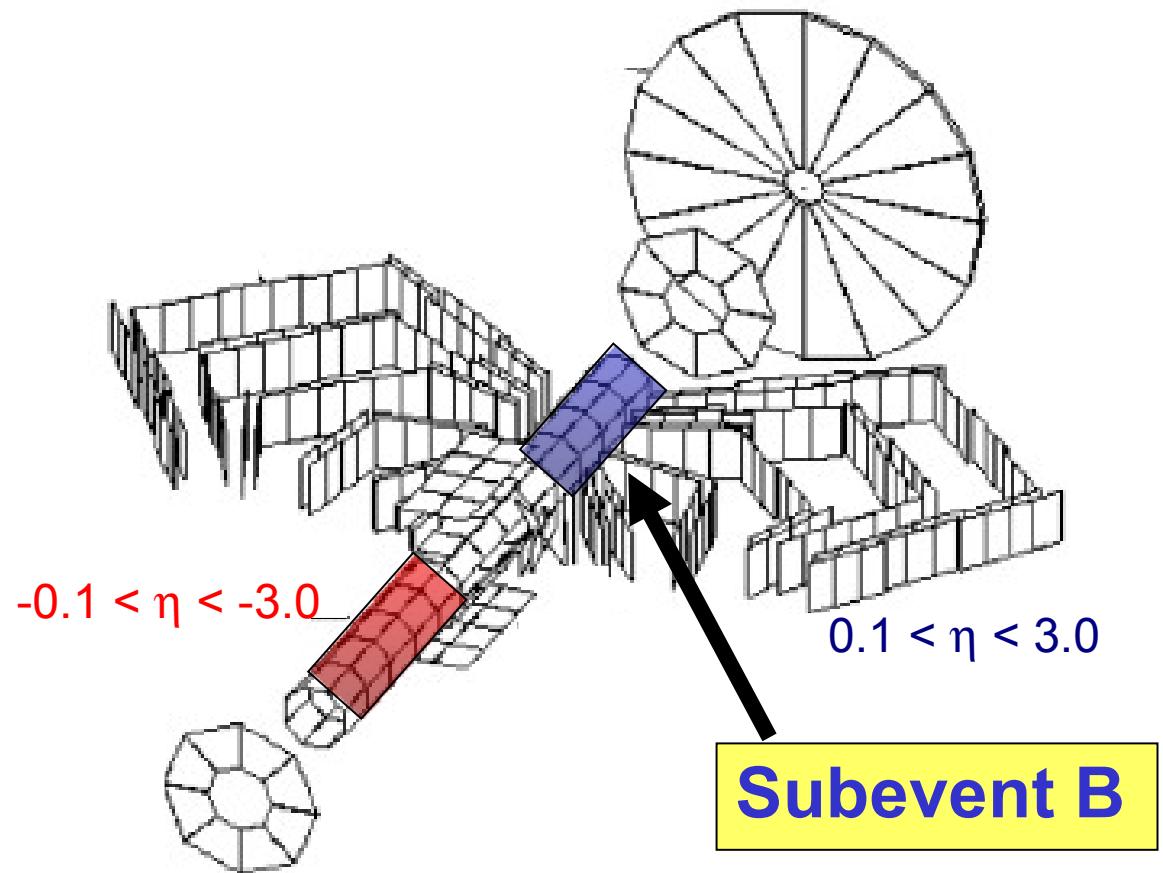
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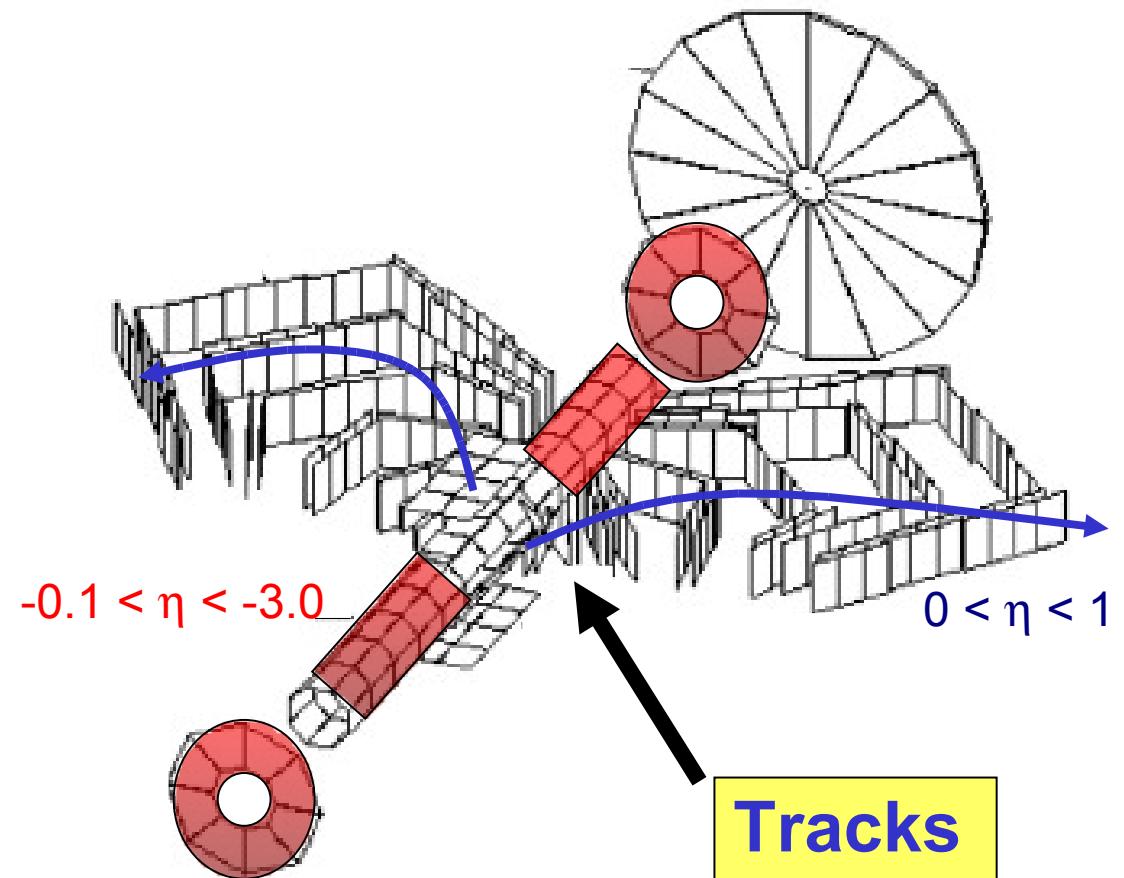
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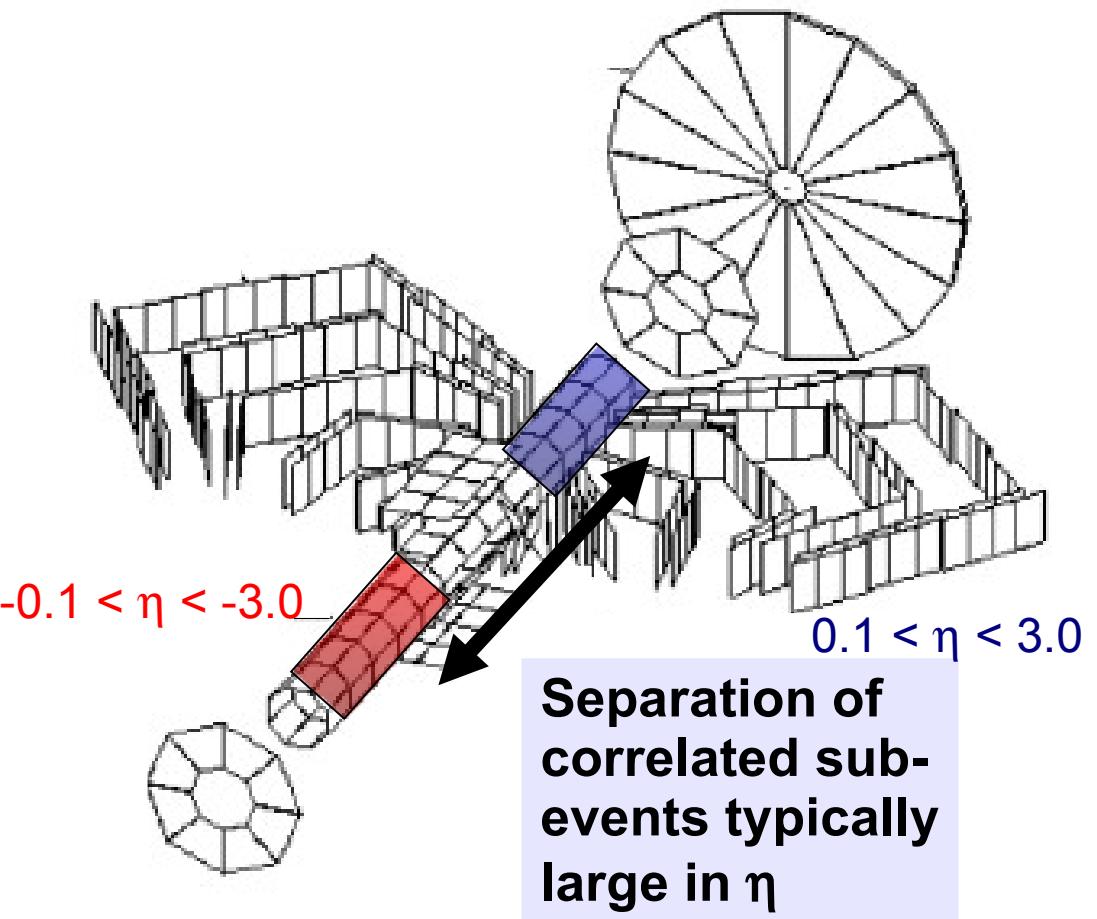
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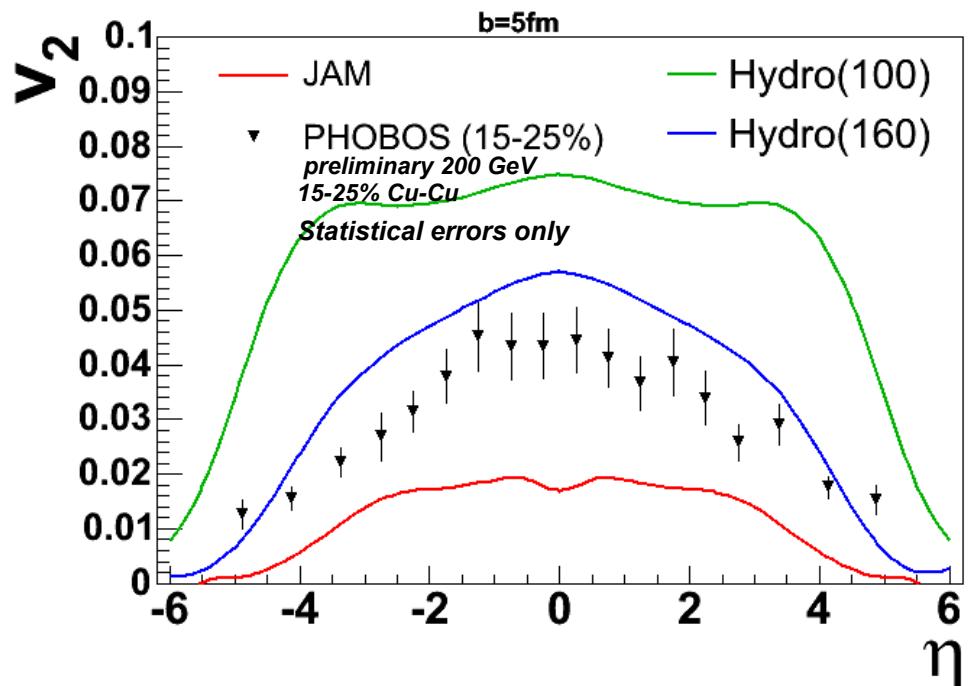
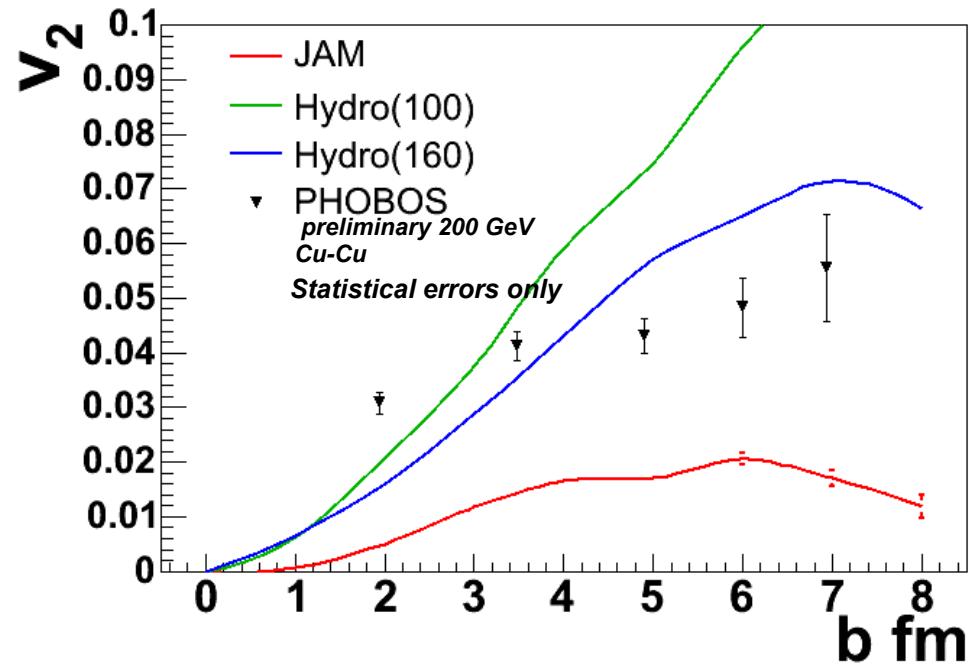
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# Elliptic flow – Cu+Cu results

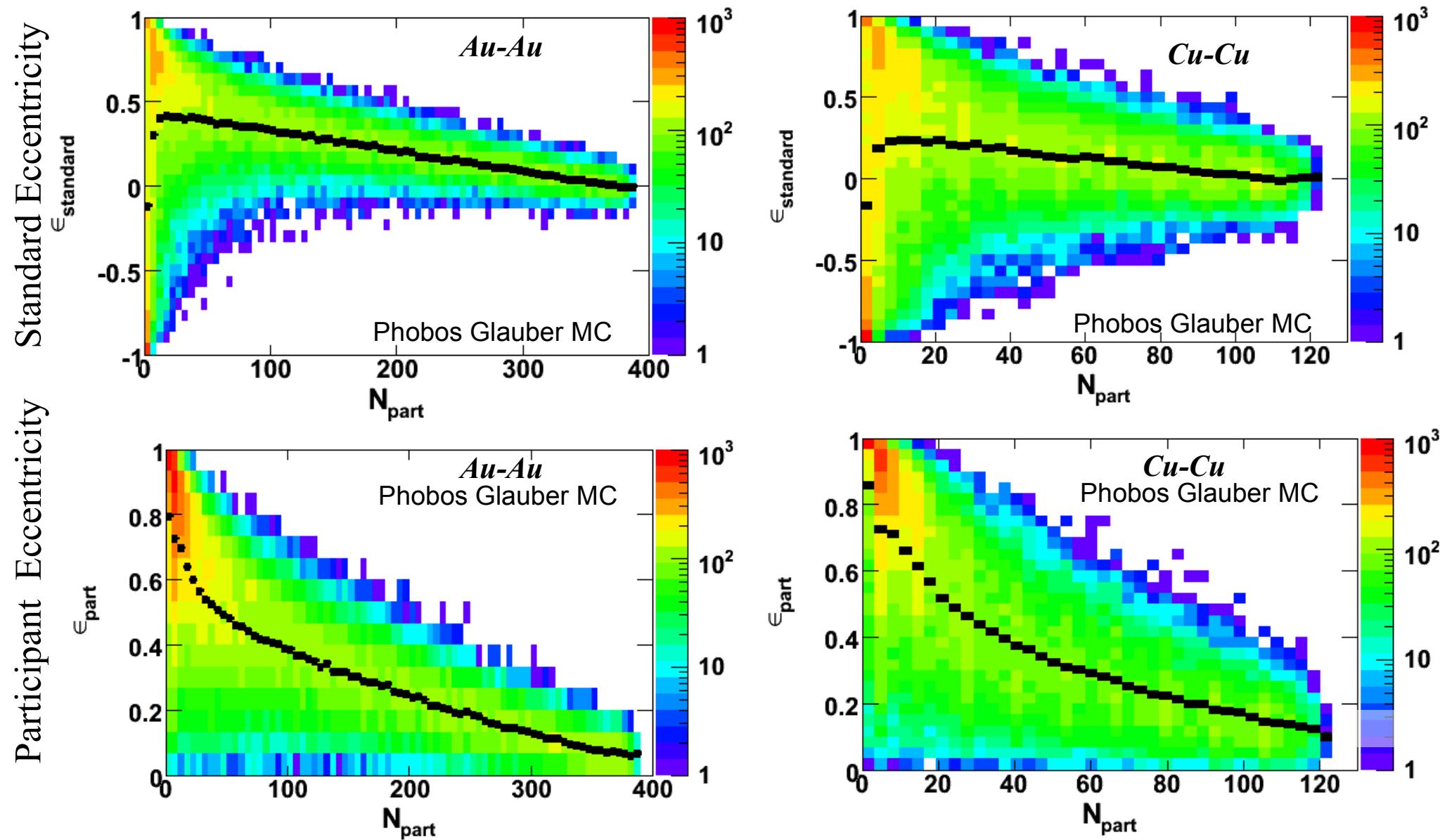


**Cu+Cu more like Hydro than JAM hadron string cascade model**

Models from Hirano et al., nucl-th/0506058; Here JAM uses a 1 fm/c formation time. Hydro (160) has kinetic freezeout temperature at 160 MeV

# Standard and participant eccentricity

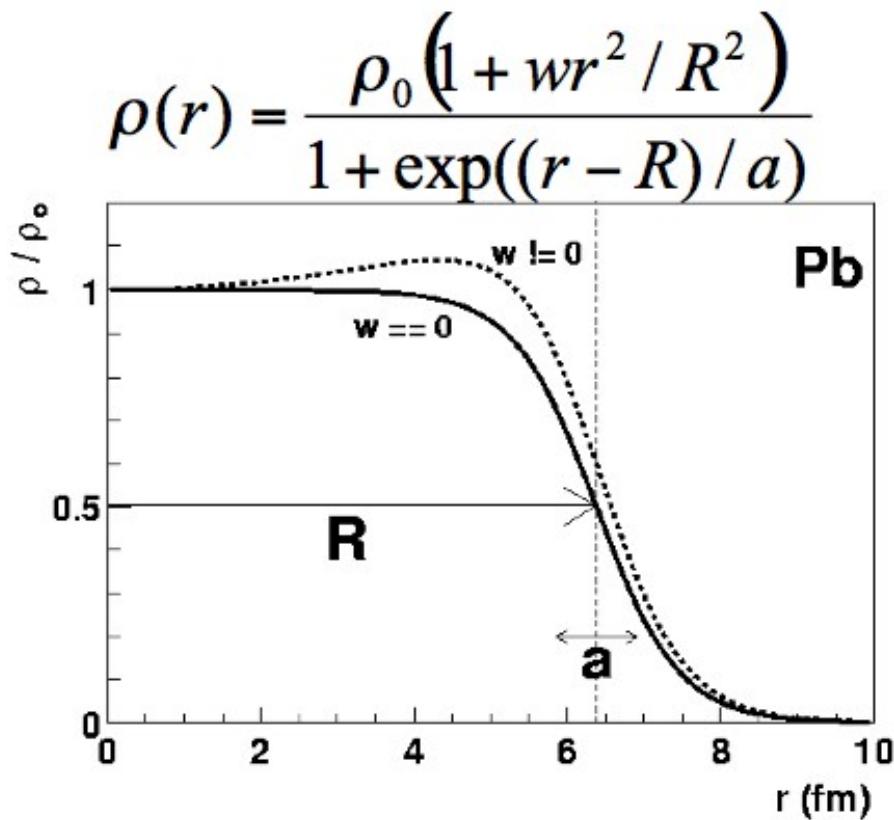
*Mean eccentricity shown in black*



Greater fluctuations in Cu+Cu. Positive fluctuations lead to non zero mean.

# Glauber Parameters Changed

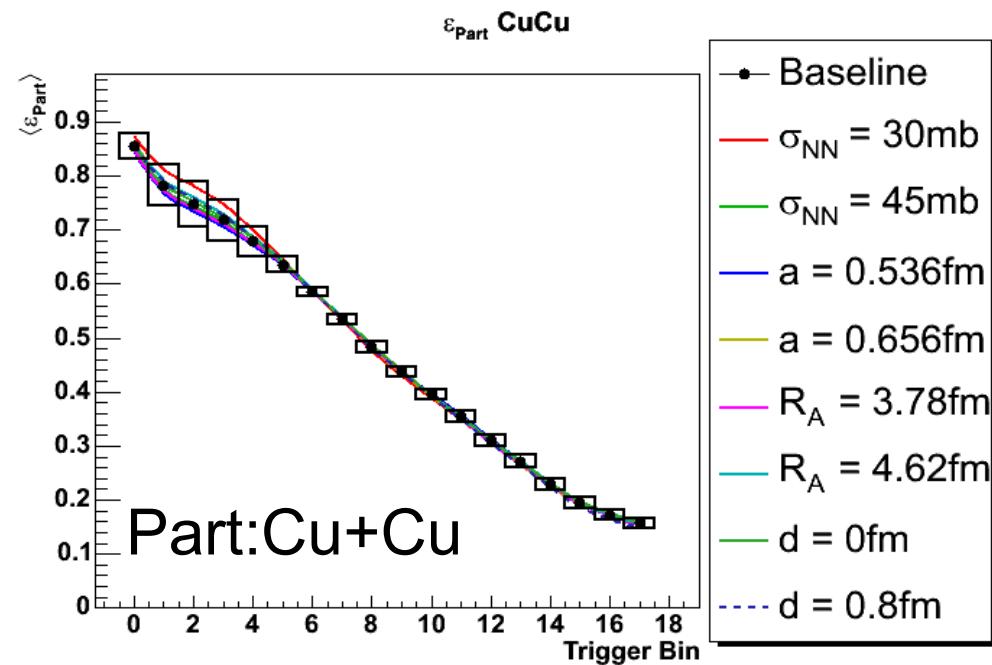
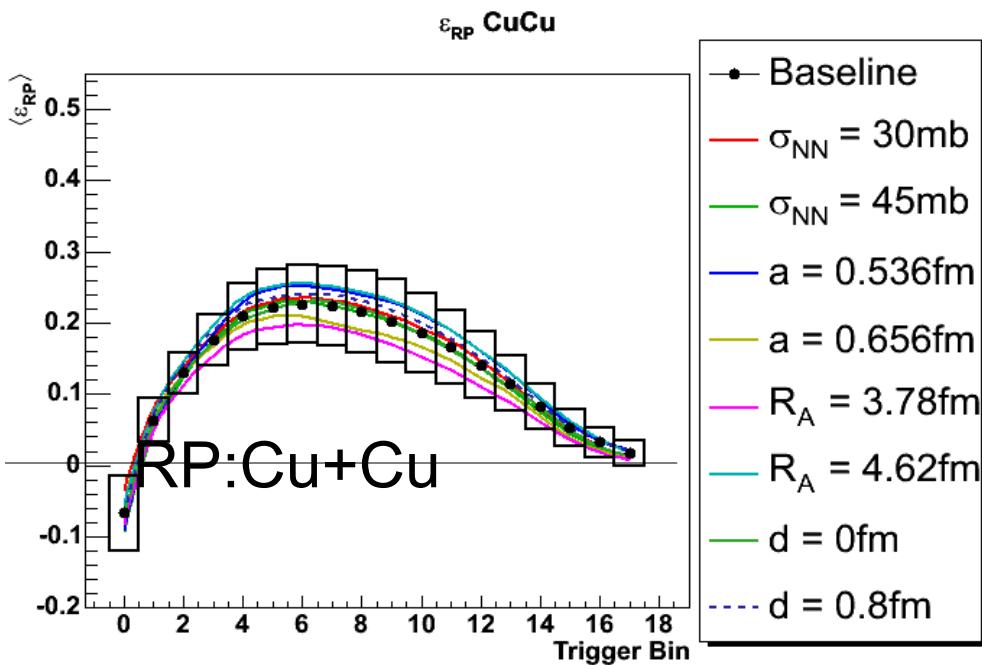
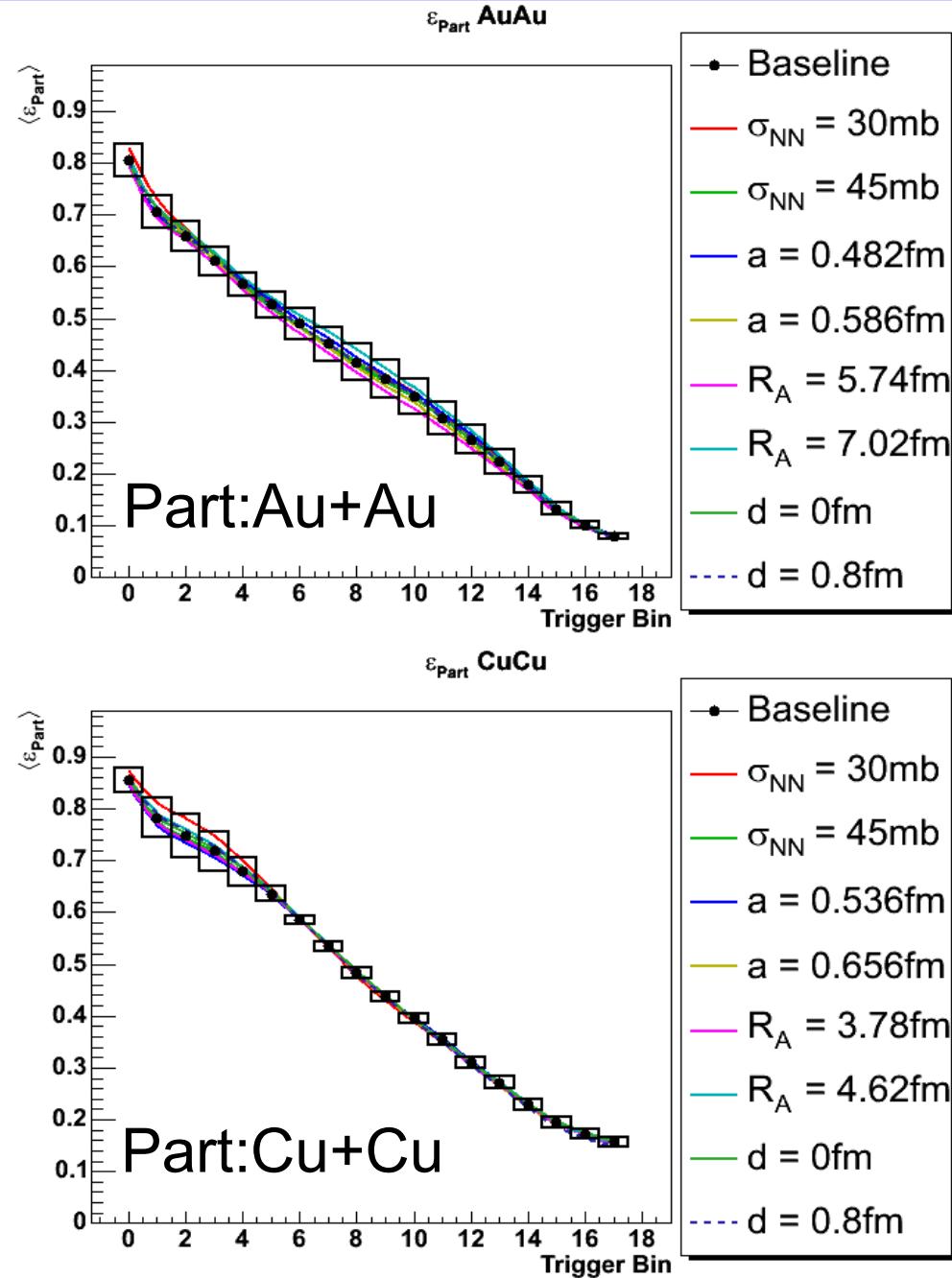
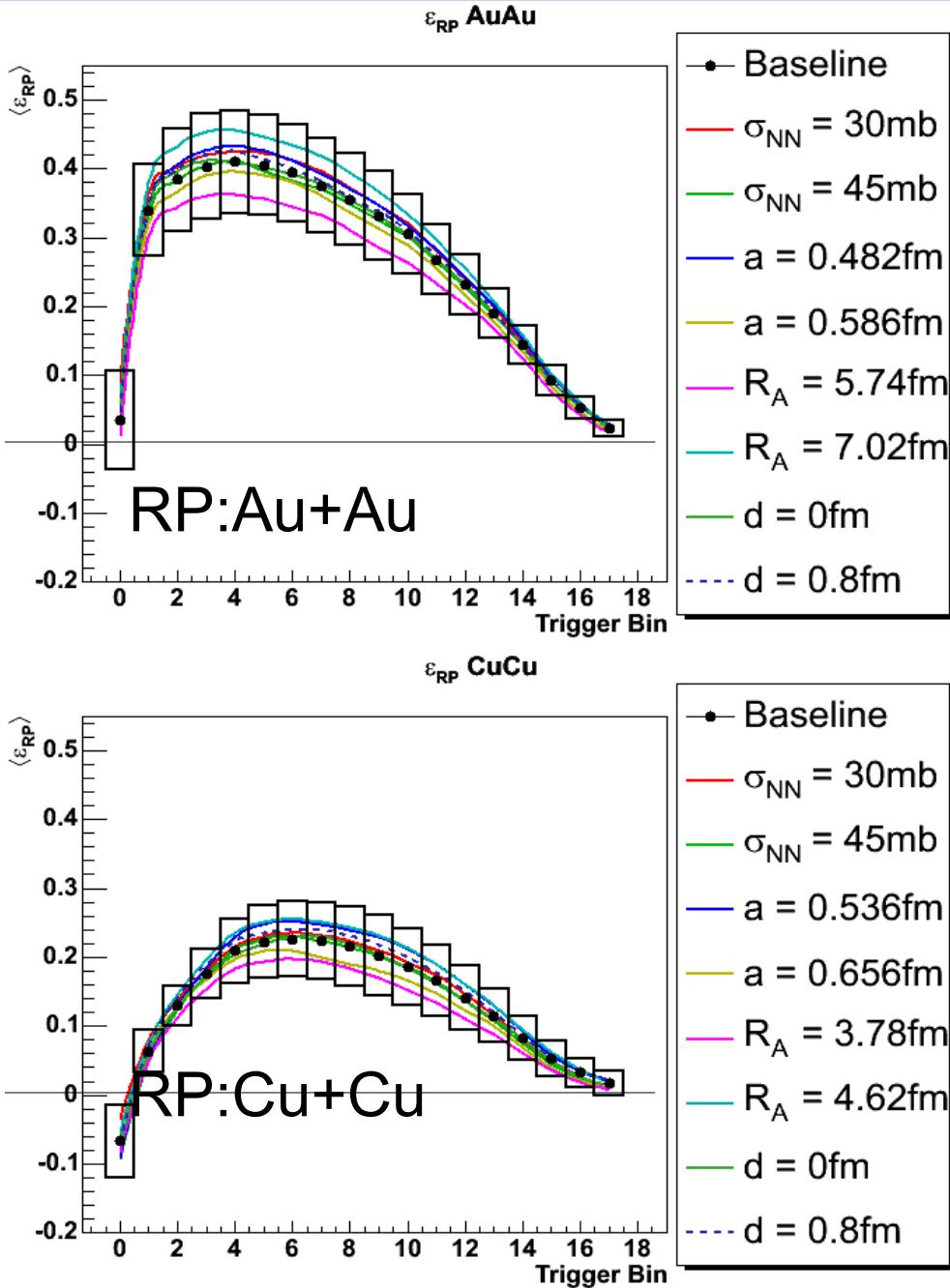
Systematic Source	Standard	How Much We Vary
Nucleon-nucleon cross-section	42 mb (for 200GeV)	30 mb (<20GeV) 45 mb (>200GeV)
Nuclear skin depth	0.535fm(Au) 0.596fm(Cu)	±10%
Nuclear radius	6.38fm (Au) 4.2fm (Cu)	±10%
Minimum nucleon separation (center-to-center)	0.4fm (like HIJING)	0fm 0.8fm



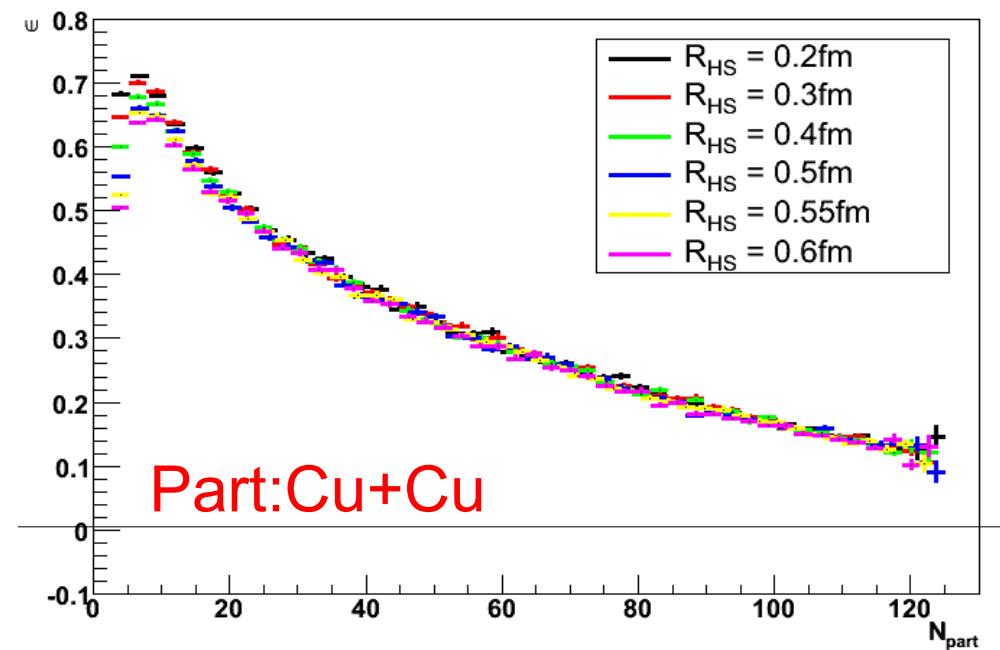
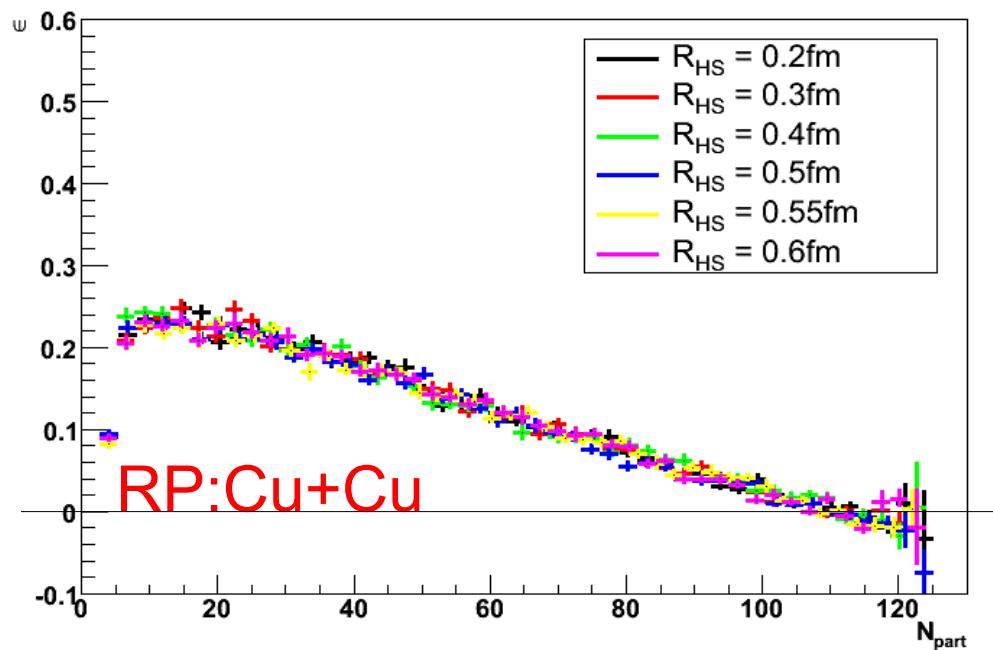
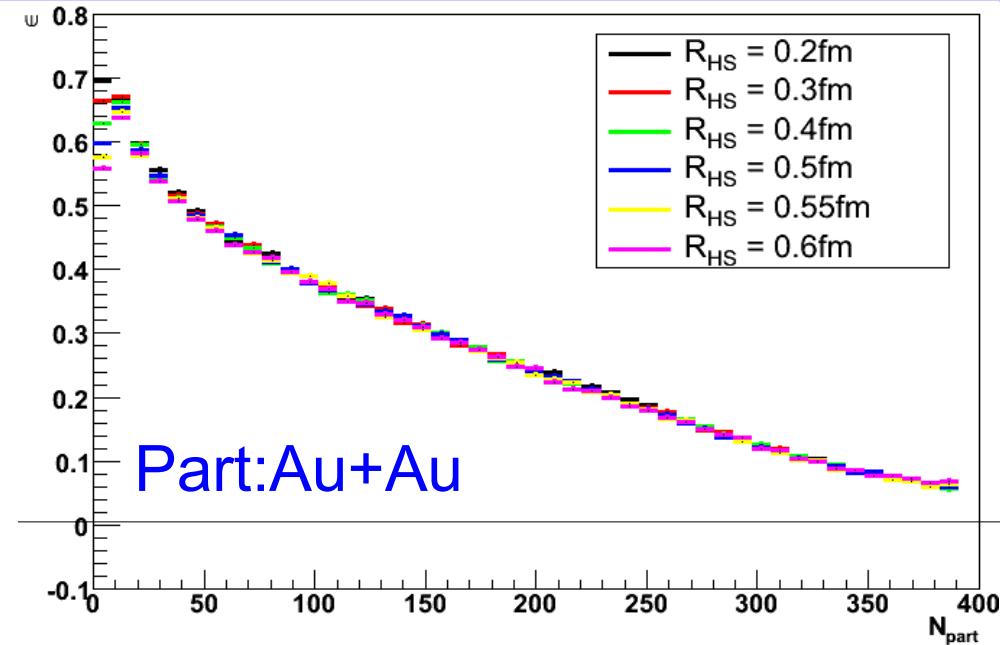
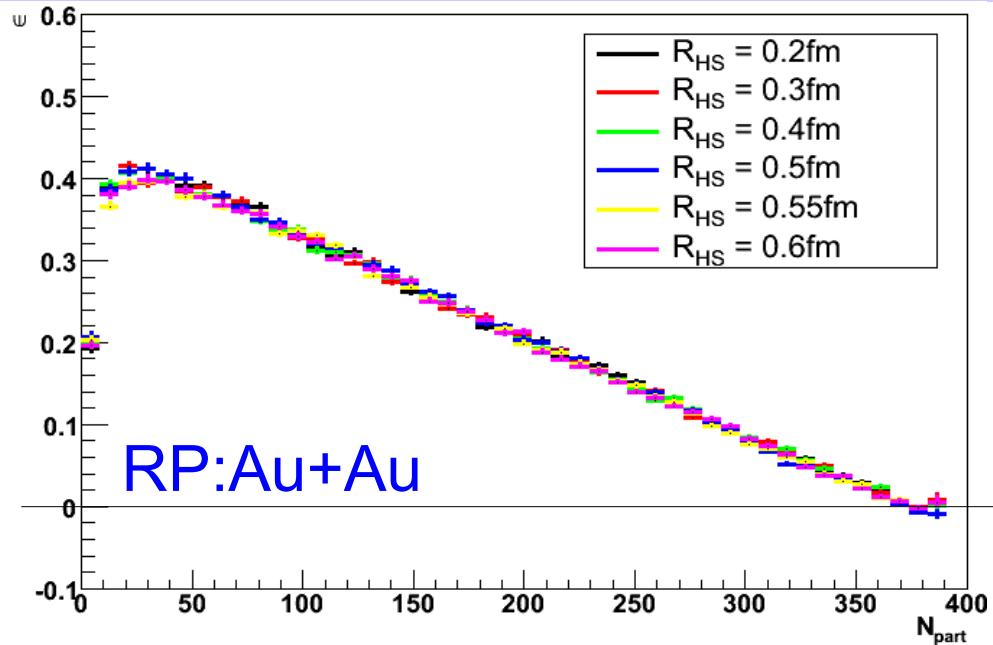
Nucleus	A	R	a	w
C	12	2.47	0	0
O	16	2.608	0.513	-0.051
Al	27	3.07	0.519	0
S	32	3.458	0.61	0
Ca	40	3.76	0.586	-0.161
Ni	58	4.309	0.516	-0.1308
Cu	63	4.2	0.596	0
W	186	6.51	0.535	0
Au	197	6.38	0.535	0
Pb	208	6.68	0.546	0
U	238	6.68	0.6	0

H. DeVries, C.W. De Jager, C. DeVries, 1987

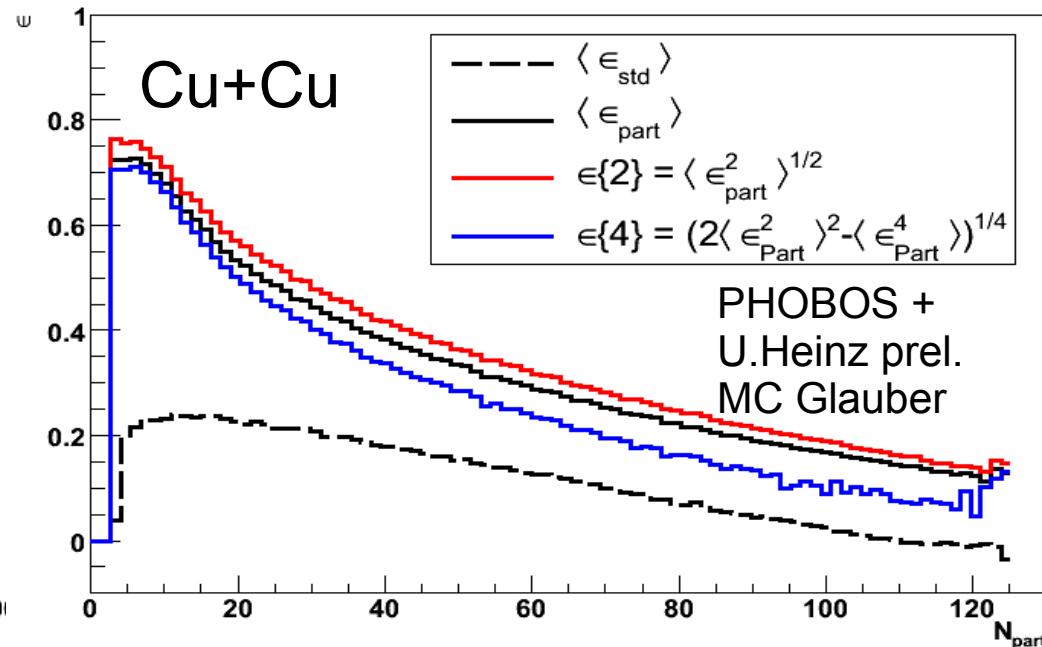
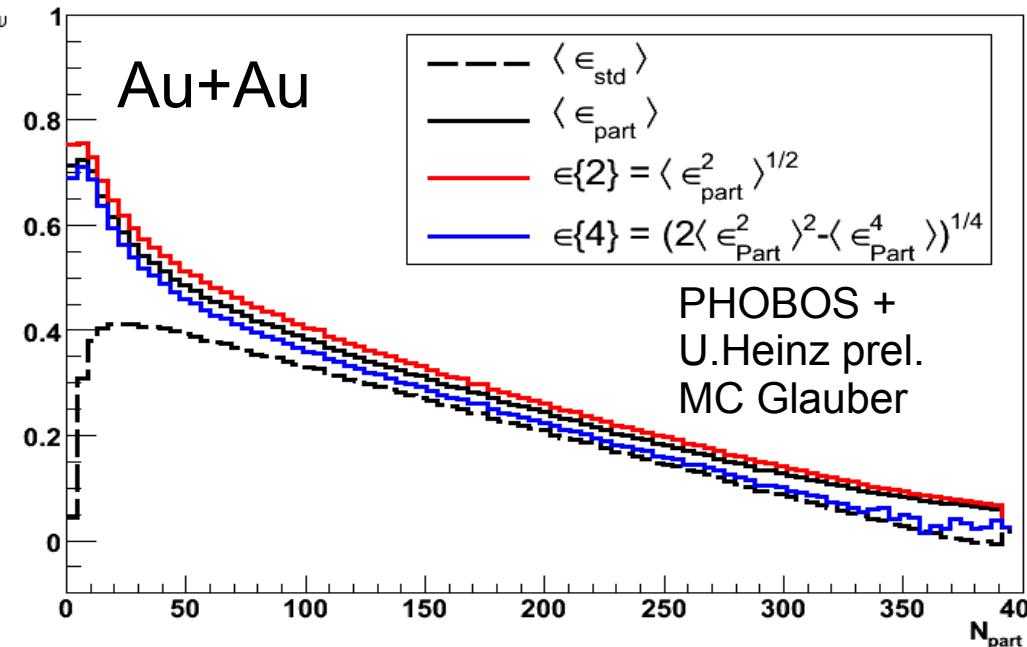
# Robustness with geometry variables



# Hard-sphere smearing (Npart weighted)



# Moments of eccentricity

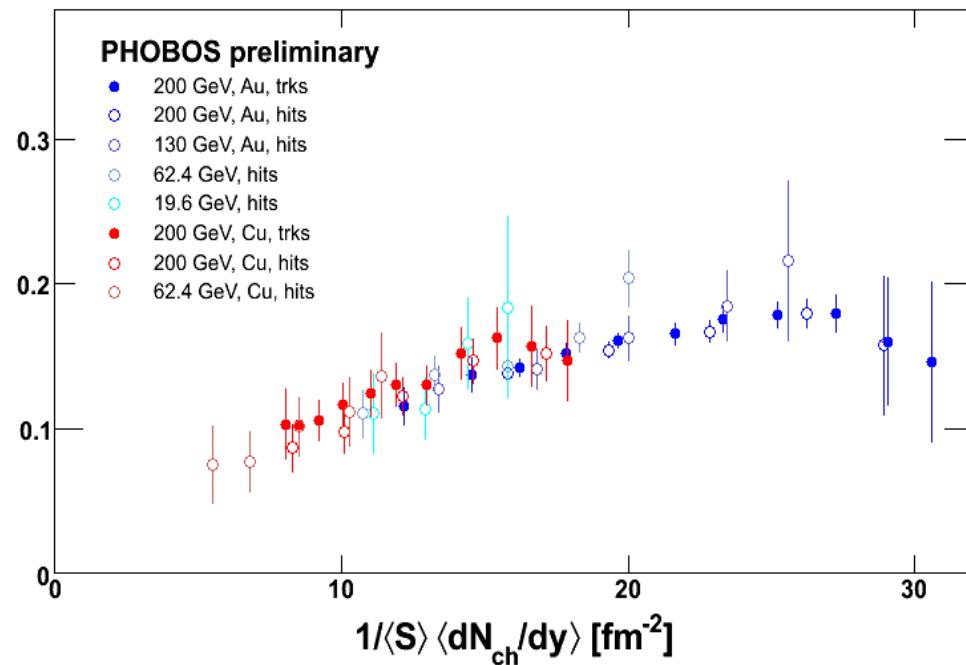
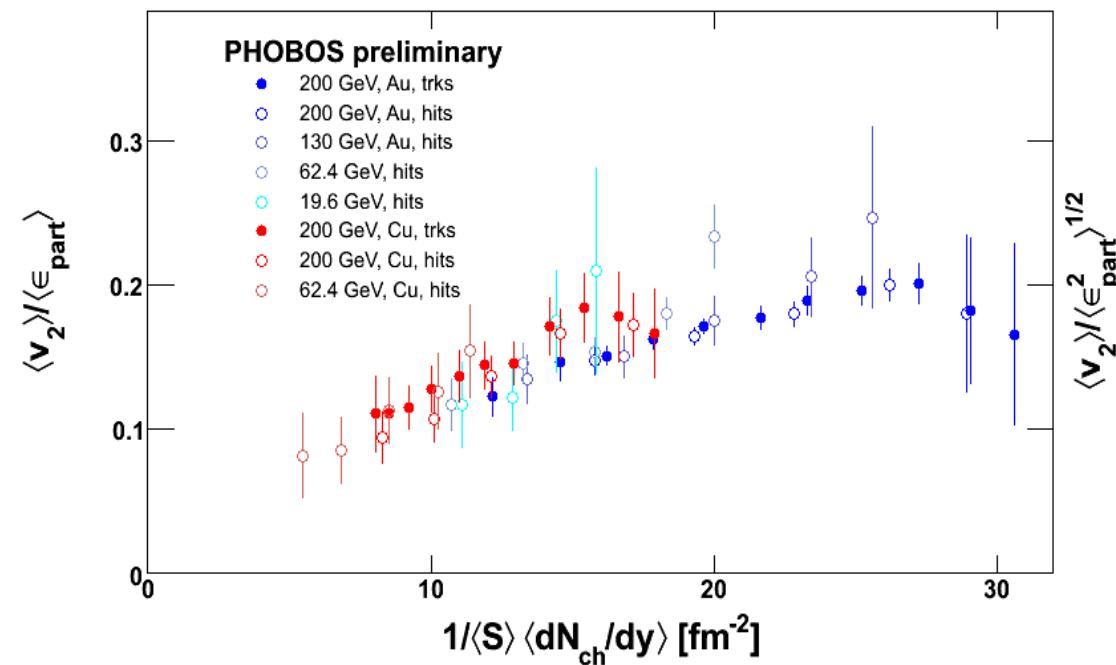


- If  $v_2$  fluctuates proportional to  $\epsilon_{\text{part}}$ , then use cumulants for  $v_2/\epsilon$  ratio
  - $\epsilon\{2\}$  rather similar to  $\epsilon_{\text{part}}$ , unlike  $\epsilon\{4\}$  that is between  $\epsilon_{\text{std}}$  and  $\epsilon_{\text{part}}$
- Outstanding question:  
To what extent is PHOBOS  $v_2\{\text{RP}\}$  sensitive to higher moments in  $v_2$ 
  - Answer probably dependent on system

M.Miller, R.Snellings, nucl-ex/0312008  
R.Bhalerao, J.Ollitrault, nucl-th/0607009

# Voloshin plots wrt to participant eccentricity

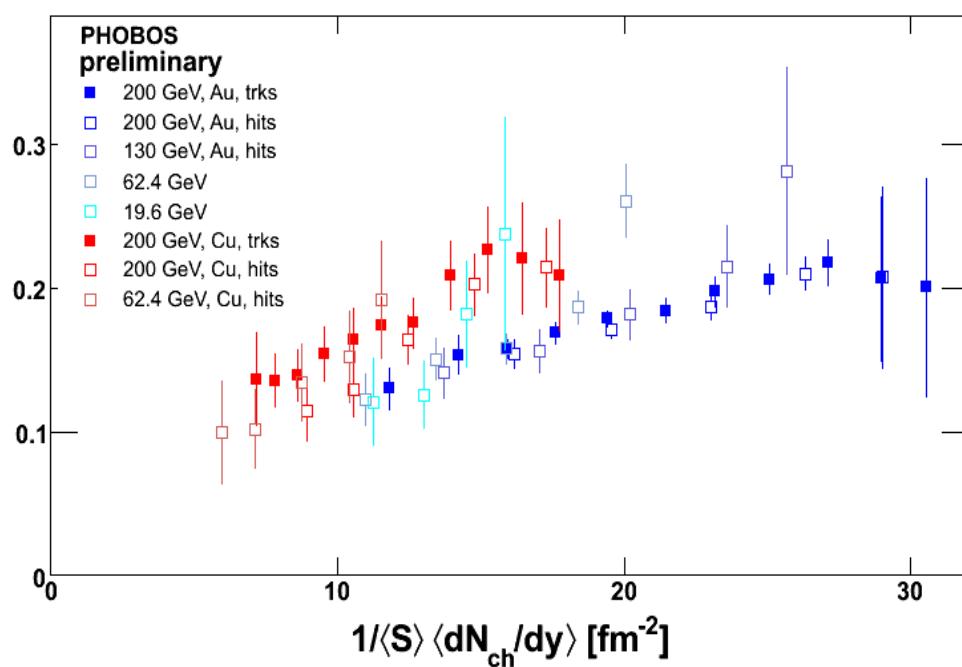
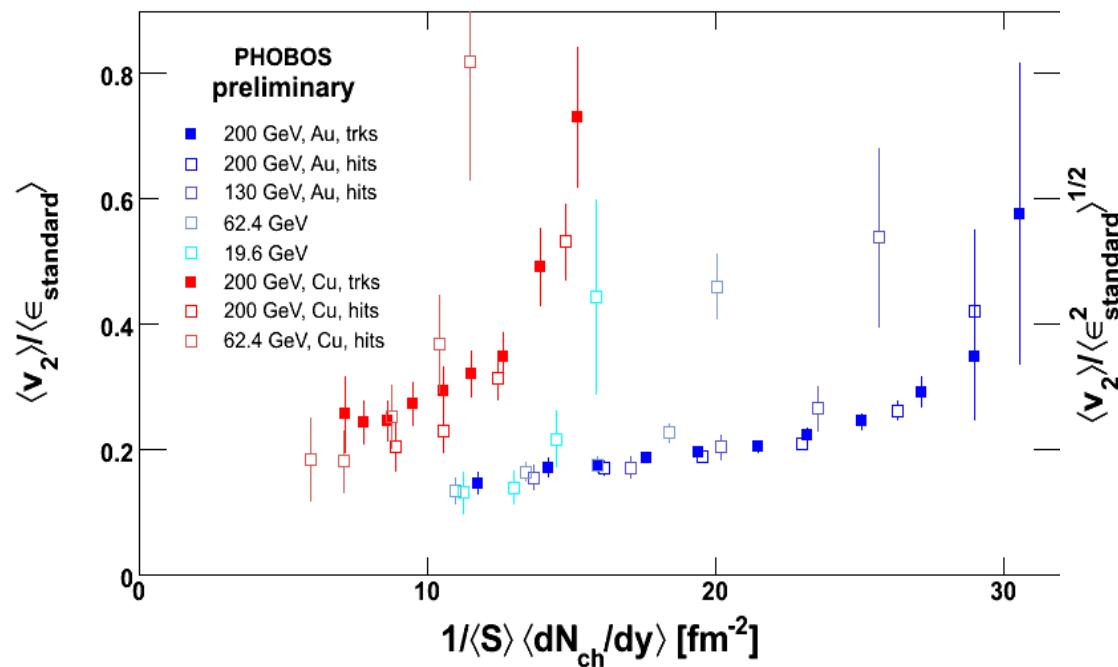
- Open question (issue):  $v_2\{\text{RP}\} \approx v_2\{2\}$  (claimed by Star, Bhalerao/Ollitrault)



- Quantitative “improvement” of the scaling

# Voloshin plots wrt to reaction-plane eccentricity

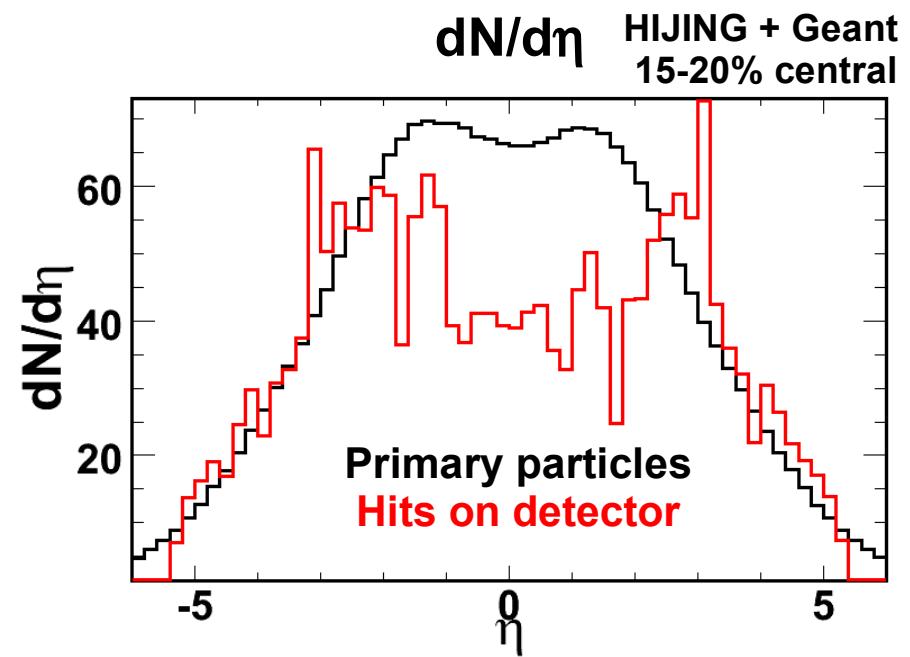
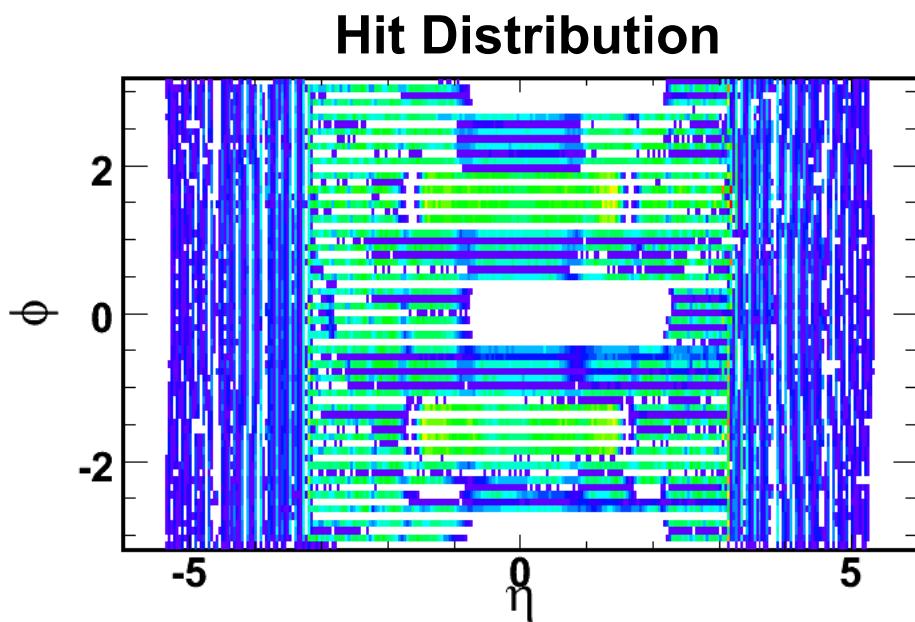
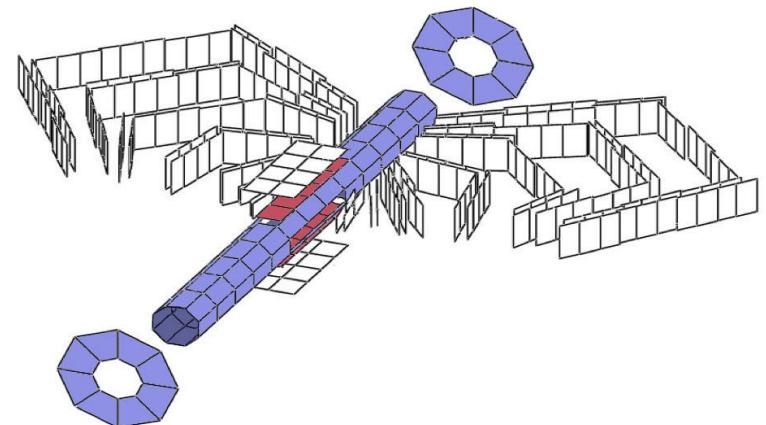
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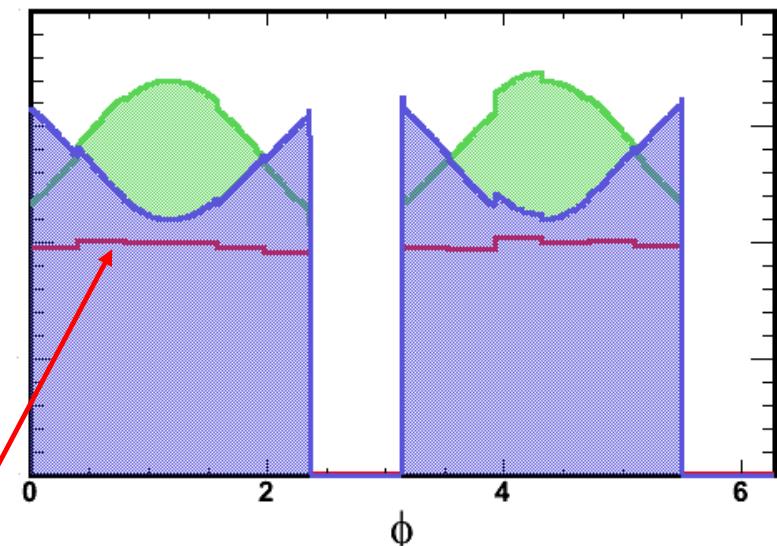
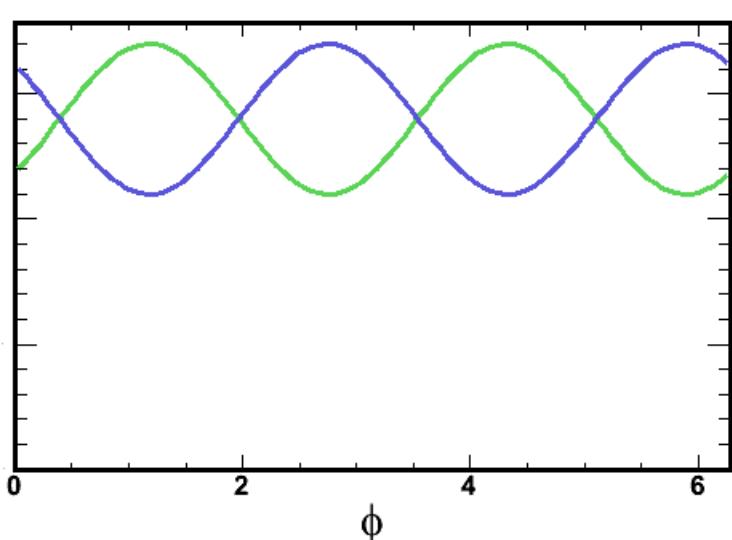
- Qualitative “improvement” of the scaling ?

# Event-by-event measurement of $v_2^{\text{obs}}$

- PHOBOS Multiplicity Array
  - $-5.4 < \eta < 5.4$  coverage
  - Holes / granularity differences
- Usage of all available information in event to determine single value for  $v_2^{\text{obs}}$



# Likelihood fit normalization



Acceptance

$$s(u, \phi_0 | \eta) = \int A(\eta, \phi) [1 + 2 u (1 - |\eta|/6) \cos(2(\phi - \phi_0))] d\phi d\eta$$

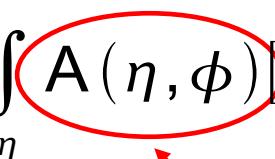
$$L(u, \phi_0) = \prod_{i=1}^n \frac{1}{s(u, \phi_0 | \eta_i)} [1 + 2 u (1 - |\eta_i|/6) \cos(2(\phi_i - \phi_0))]$$

# Event-by-event measurement of $v_2^{\text{obs}}$

$$L(v_2^{\text{obs}}, \phi_0) = \prod_{i=1}^n p(\eta_i) [1 + 2v_2^{\text{obs}} (1 - \frac{|\eta_i|}{6}) \cos(2\phi - 2\phi_0)]$$

- Maximize likelihood to find “most likely” value of  $v_2^{\text{obs}}$  (and  $\phi_0$ )
- Comparing values of  $v_2^{\text{obs}}$  and  $\phi_0$ 
  - PDF folded by acceptance must be normalized to the same value for different  $v_2^{\text{obs}}$  and  $\phi_0$ 's

$$p(v_2^{\text{obs}}, \phi_0; \eta) = 1 / \int_{\eta} A(\eta, \phi) [1 + 2v_2^{\text{obs}} (1 - \frac{|\eta|}{6}) \cos(2\phi - 2\phi_0)] d\eta d\phi$$

  
Acceptance

# Determining the kernel

- In a single bin of  $v_2$  and  $n$

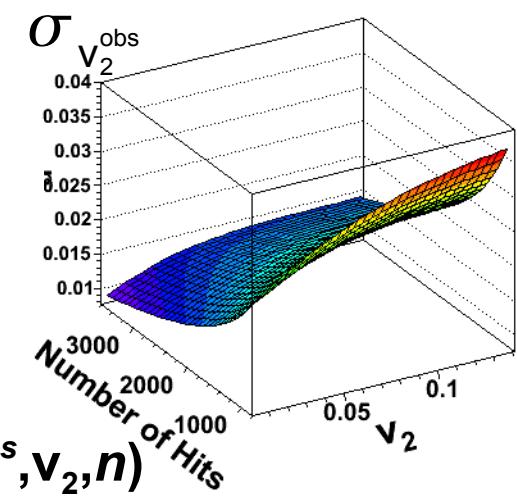
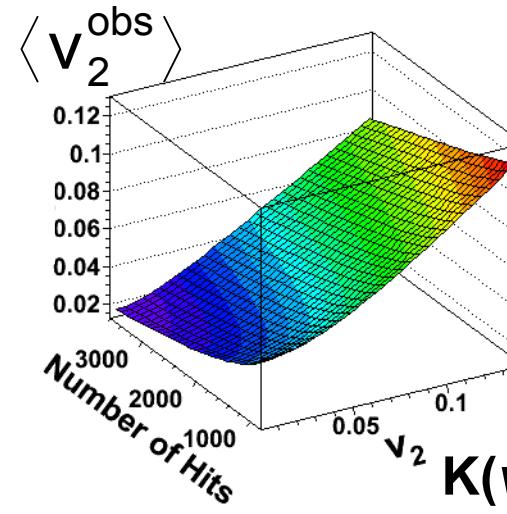
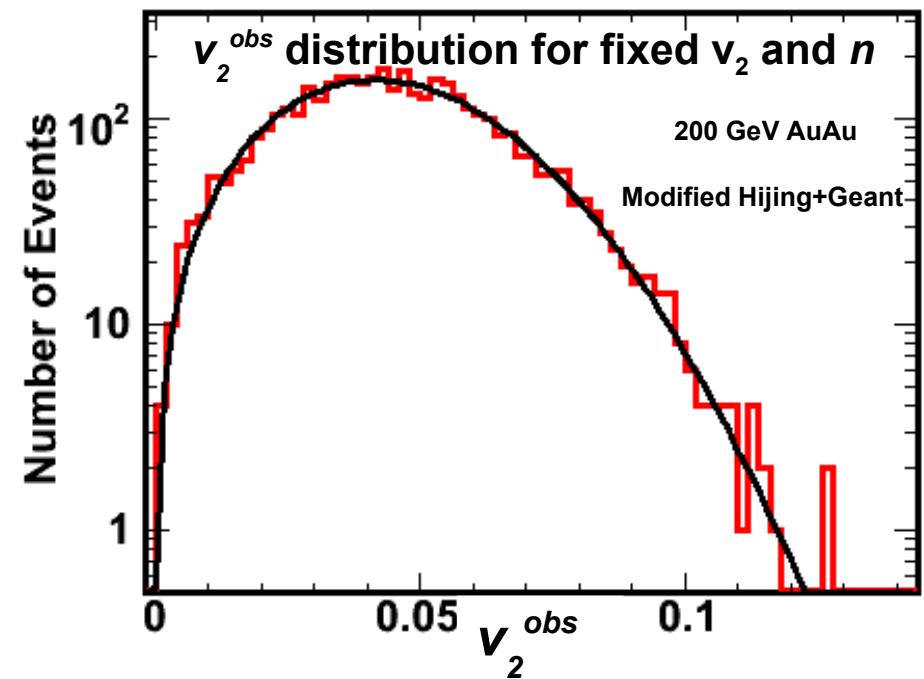
$$K(v_2^{\text{obs}}; v_2, n) = v_2^{\text{obs}} \exp\left(\frac{-(v_2^{\text{obs}} - a)^2}{2b^2}\right)$$

- Correspondance

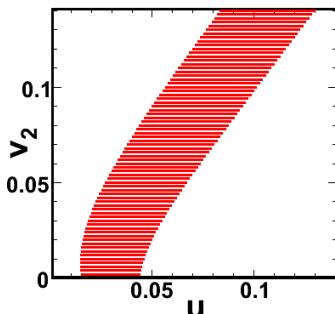
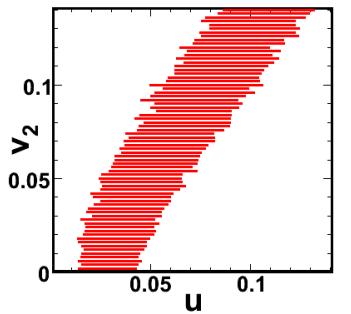
$$(a, b) \Leftrightarrow (\langle v_2^{\text{obs}} \rangle, \sigma_{v_2^{\text{obs}}})$$

- Determine  $\langle v_2^{\text{obs}} \rangle$  and  $\sigma_{v_2^{\text{obs}}}$  in bins of  $v_2$  and  $n$  and determine  $(a, b)$
- Fit smooth functions
- Integrate out multiplicity

$$K(v_2^{\text{obs}}; v_2) = \int K(v_2^{\text{obs}}; v_2, n) N(n) dn$$



# Calculating the kernel: Functions observed to fit

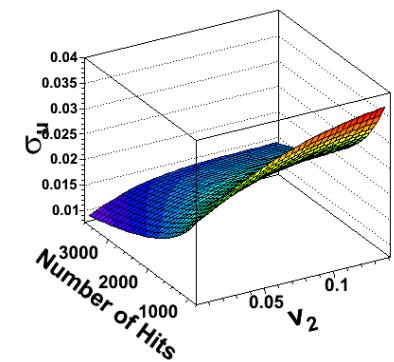
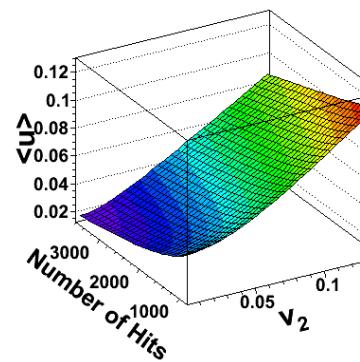
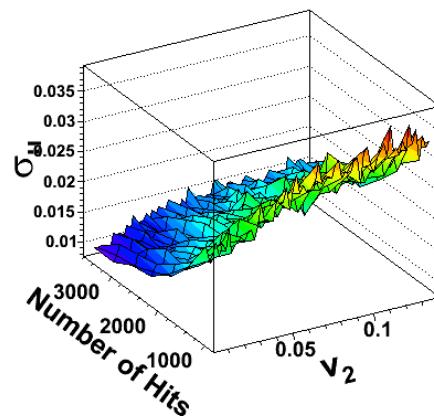
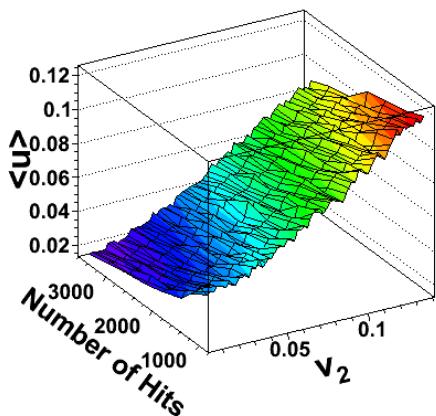


$$\langle u \rangle(v_2)|_n = \sqrt{m_1 \cdot v_2^2 + m_2}$$

$$\sigma_u(v_2)|_n = \frac{1}{r_1 + \frac{2}{3} e^{r_2 v_2}}$$

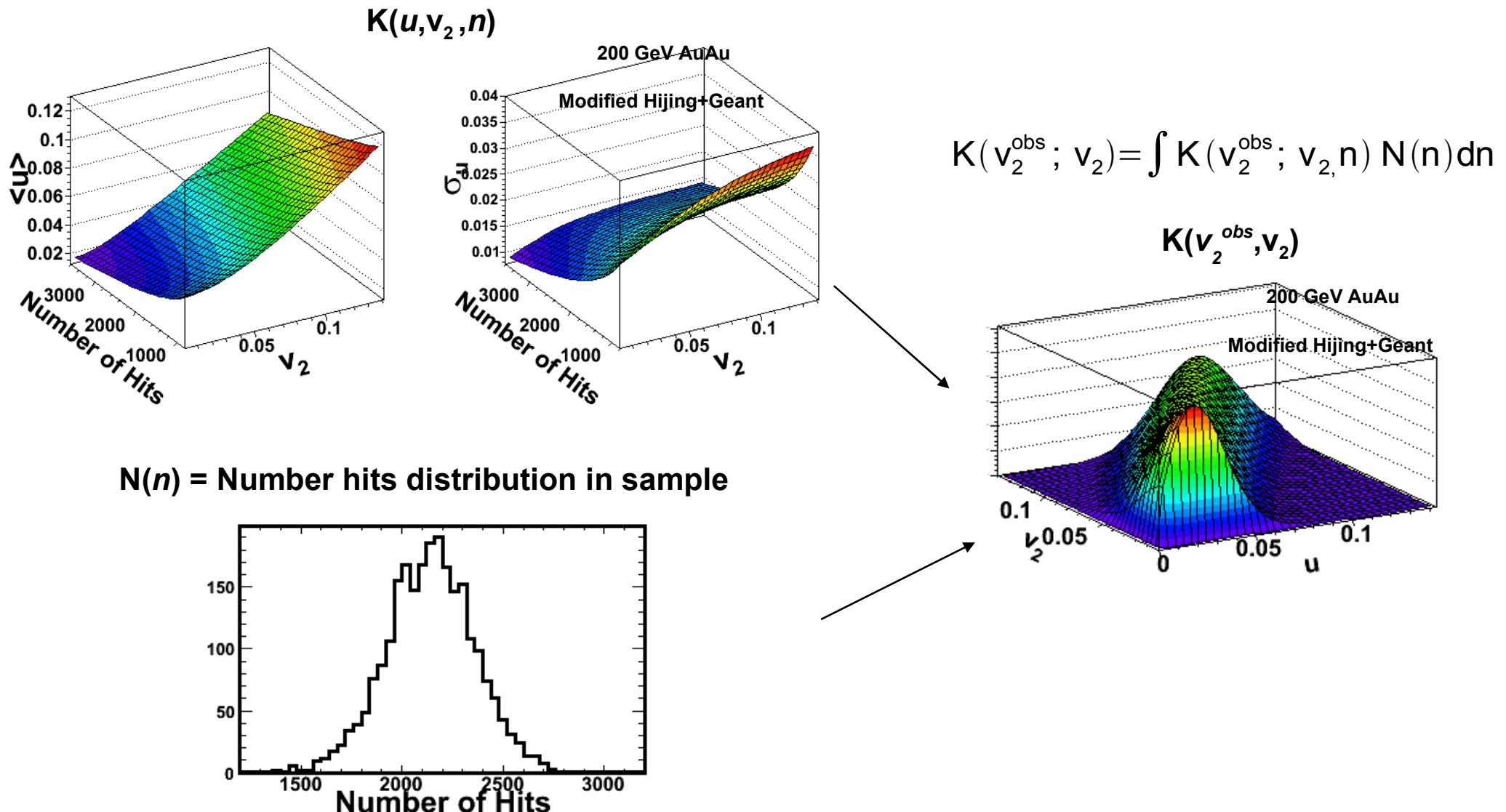
$$\langle u \rangle(v_2, n) = \sqrt{(M_1 n^2 + M_2 n + M_3) \cdot v_2^2 + (M_4 * n + M_5)}$$

$$\sigma_u(v_2, n) = \frac{R_1}{(R_2 \sqrt{n} + 1) \cdot (1 + \frac{2}{3} e^{(R_3 n + R_4) v_2})}$$



# Determining the kernel

- Multiplicity dependence can be integrated out

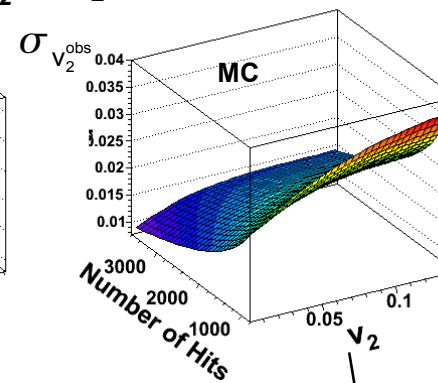
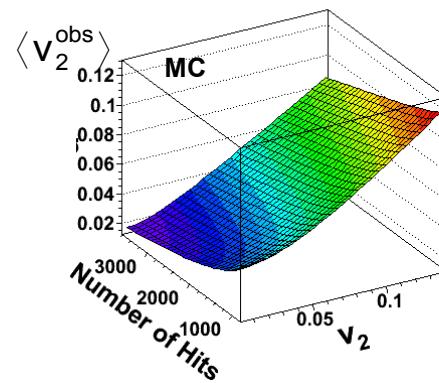


$N(n)$  = Number hits distribution in sample

# Summary of flow fluctuation measurement method

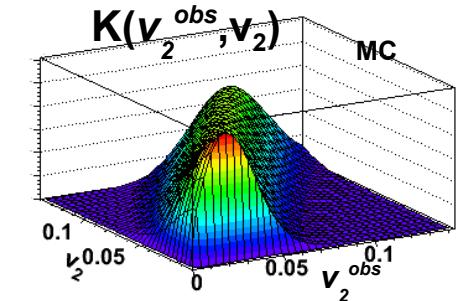
Many MC events →

$$K(v_2^{obs}, v_2, n)$$



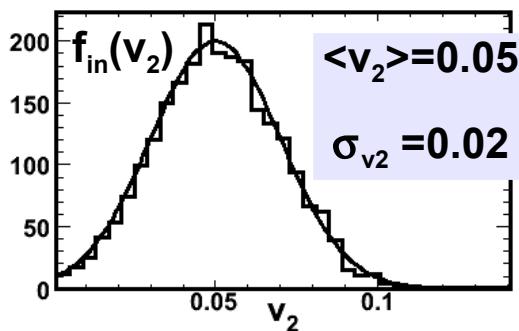
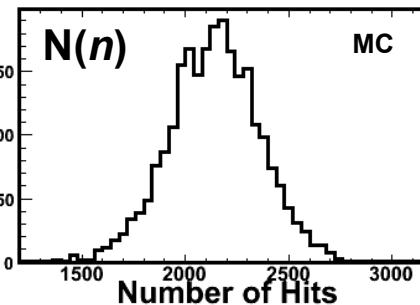
Integration

$$K(v_2^{obs}; v_2) = \int K(v_2^{obs}; v_2, n) N(n) dn$$

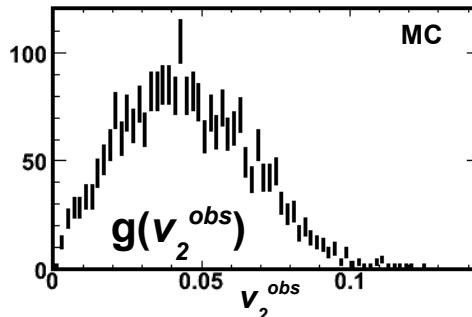


Input Sample

measurement



measurement



Minimize  $\chi^2$  in integral

$$g(v_2^{obs}) = \int_0^\infty K(v_2^{obs}, v_2) f(v_2) dv_2$$

